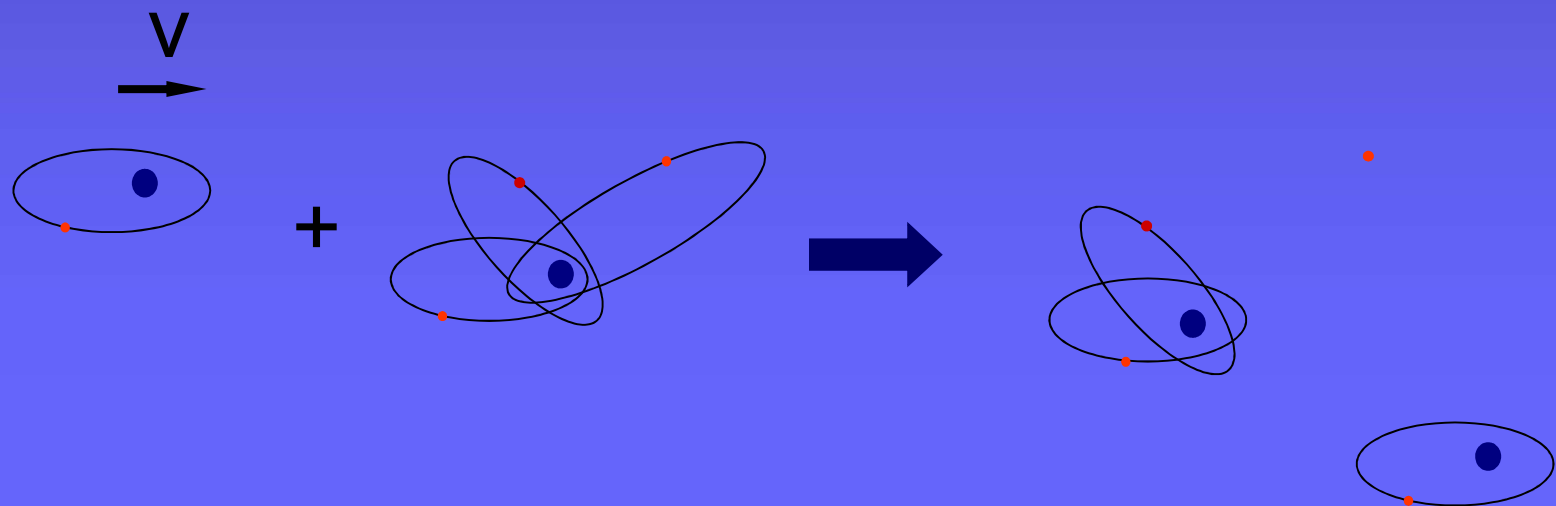


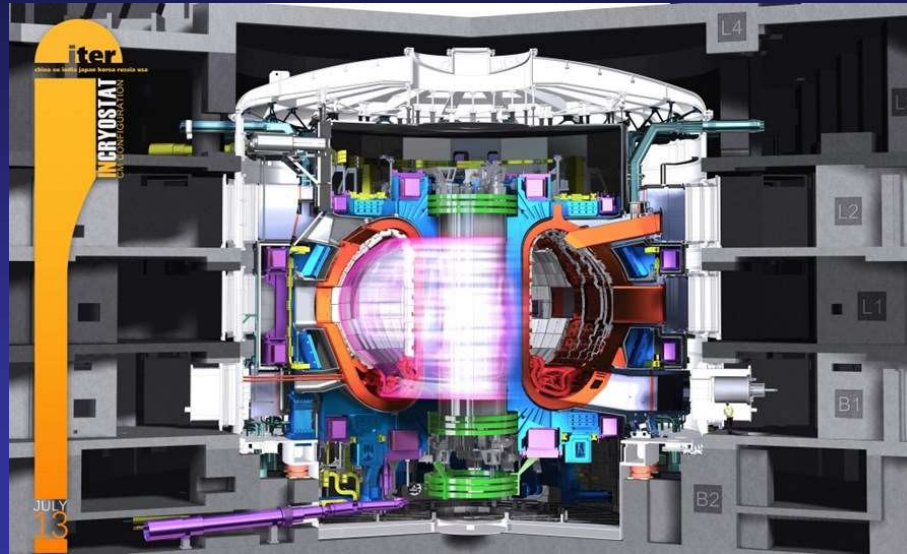
Ionization, total and state selective charge exchange cross sections in fusion related collision systems

Károly Tőkési

Institute for Nuclear Research, Debrecen, Hungary

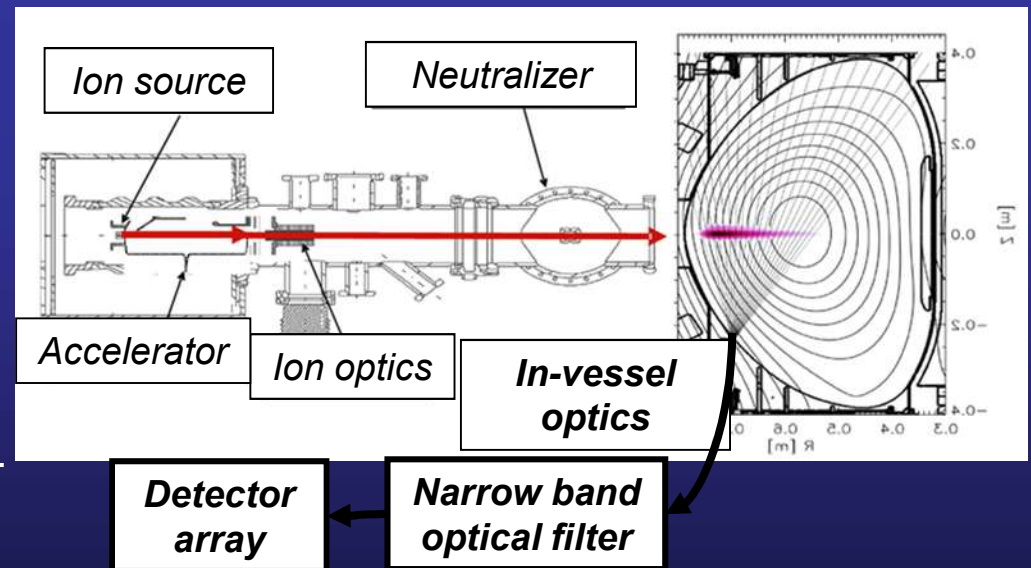


The aim



BES diagnostics

- Active plasma diagnostics procedure
- Use of H-type of atoms such as D, Li, Na. (which possess **one valence electron**)
 - Heating beams (H, D)
 - Diagnostic beams (Li, Na)
- Purpose: **density** and **fluctuation** meas.
 - Fluct. timescale: 10 – 200kHz
 - Fluct. spatial scale: 1 – 4 cm



Outlook

Basic idea – Classical treatments

Theory

Classical Trajectory Monte Carlo (CTMC) model

Quasi-Classical Trajectory Monte Carlo (QCTMC) model

Results

1. Data for Atomic Processes of Neutral Beams in Fusion Plasma

$\text{Be}^{4+} + \text{H}(1s)$ collisions

H + H type collisions

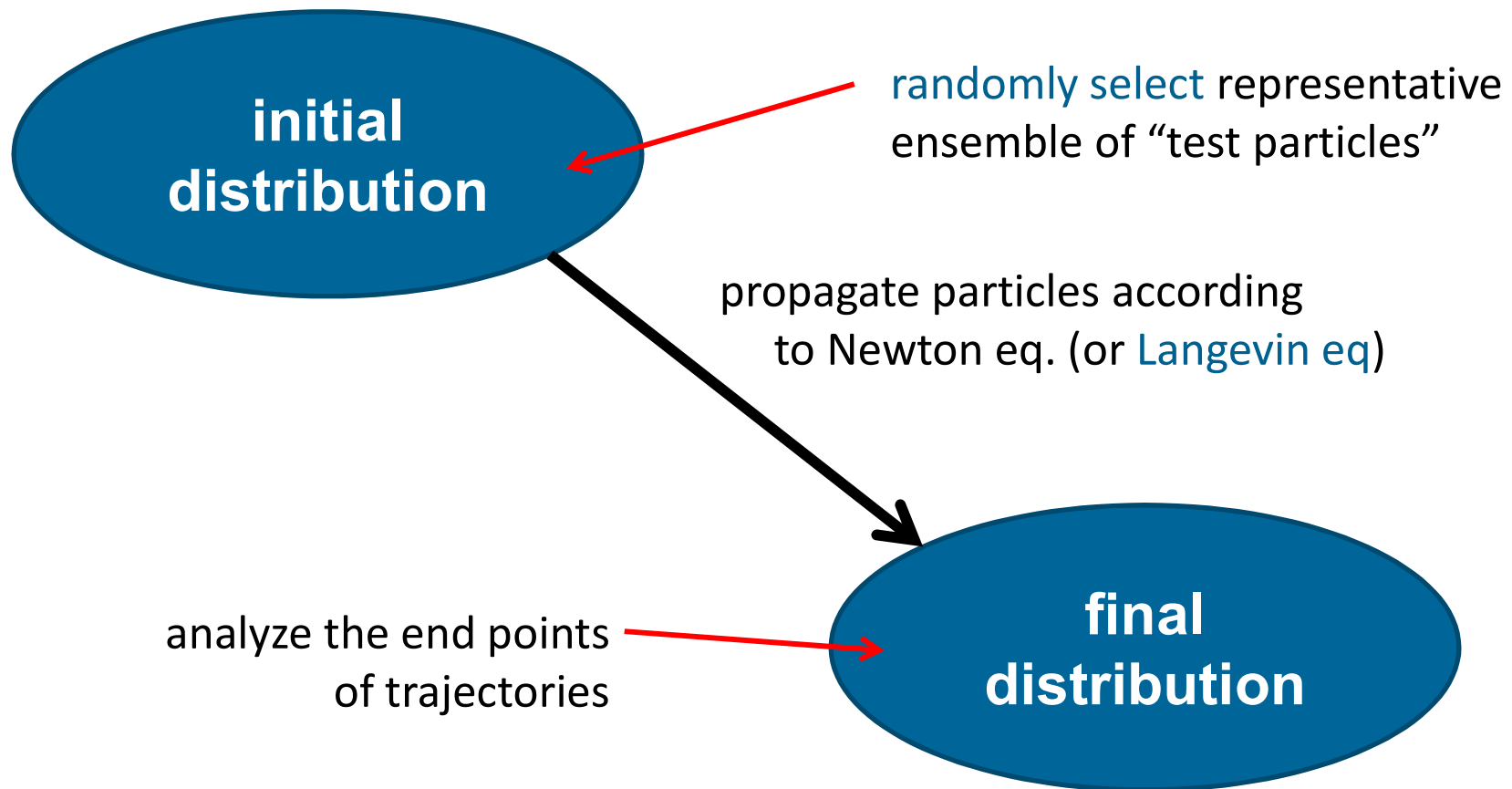
2. Injected Impurities

$\text{Li}^+ + \text{He}$, $\text{Li}^+ + \text{N}$

Summary

Approximations: CTMC simulations

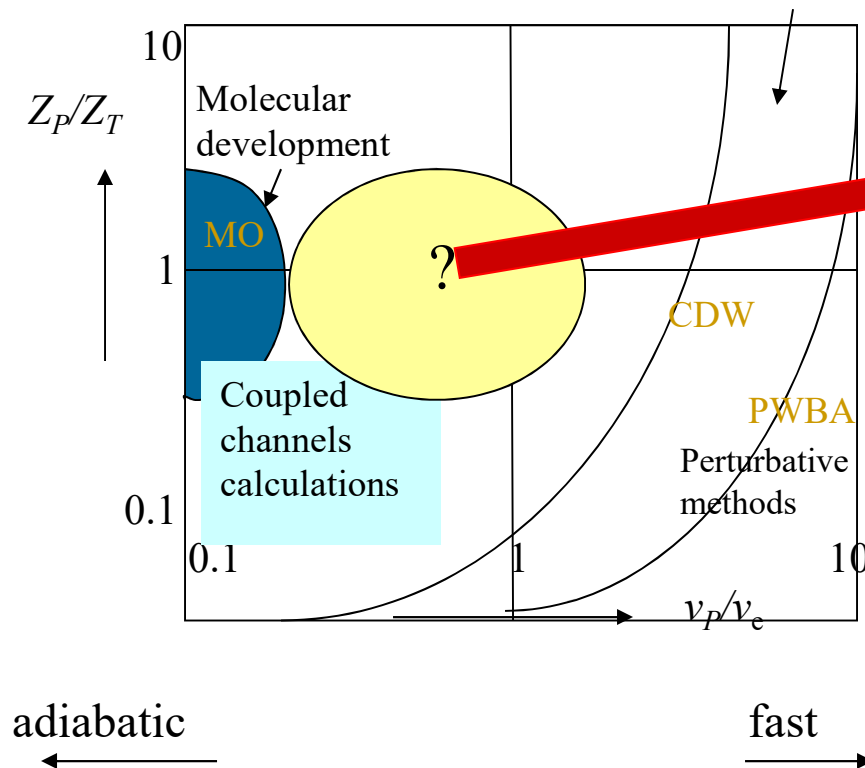
Flow diagram for a MC simulation:



Ionization in ion-atom collisions

Description:

Distorted wave approximations



Non-perturbative models:

Classical Trajectory Monte Carlo (CTMC) method

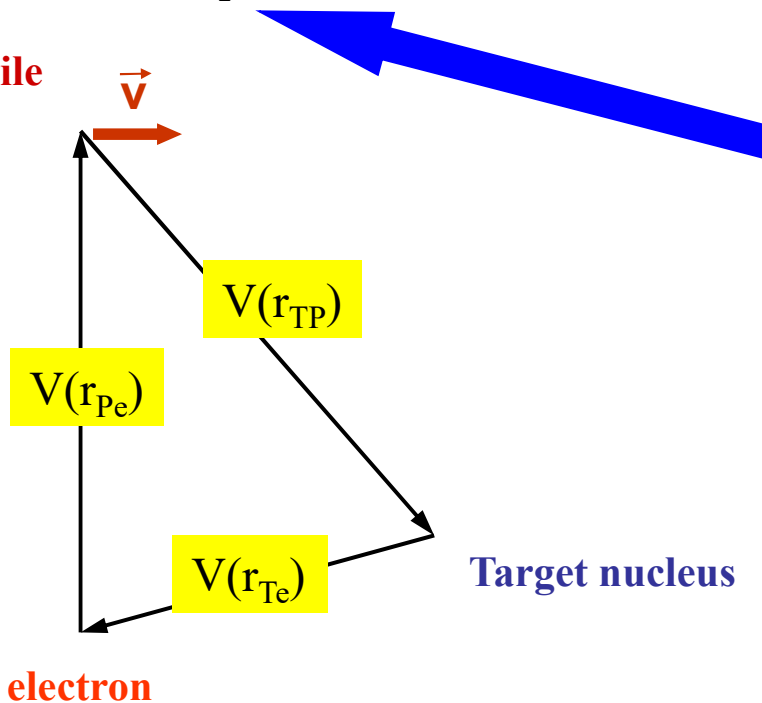
3-body CTMC approach

- **Classical nonperturbative method** – „theoretical experiment”
- **Treats the many-body interactions**

Model potential:

$$V(\mathbf{r}) = -\frac{(Z-1)\Omega(r)+1}{r}, \quad \text{where } \Omega(r) = [Hd(e^{r/d}-1)+1]^{-1}$$

Projectile



Specific for the present work:

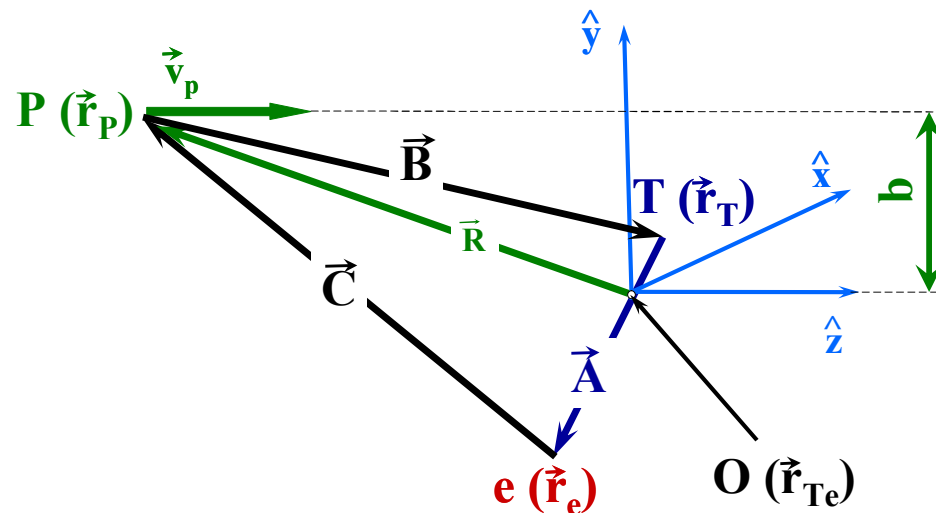
-Screened core potentials for both partners
(analytic GSZ model pot.)

-Strategies for extracting the relevant information

- a three-body balance is bound by E and \mathbf{p} conservation;
- final-state kinematics does not provide information about the mechanism

CTMC approach

- **Classical nonperturbative method**
 - „theoretical experiment”
- **Treats the many-body interactions**



$$L = L_K - L_V$$

$$L_K = \frac{1}{2} m_P \dot{\vec{r}}_P^2 + \frac{1}{2} m_e \dot{\vec{r}}_e^2 + \frac{1}{2} m_T \dot{\vec{r}}_T^2$$

$$L_V = \frac{Z_P(|\vec{r}_P - \vec{r}_e|)Z_e}{|\vec{r}_P - \vec{r}_e|} + \frac{Z_P(|\vec{r}_P - \vec{r}_T|)Z_T(|\vec{r}_P - \vec{r}_T|)}{|\vec{r}_P - \vec{r}_T|} + \frac{Z_e Z_T(|\vec{r}_e - \vec{r}_T|)}{|\vec{r}_e - \vec{r}_T|}$$

Classical principal number

$$n_c = Z_T Z_e \left(\frac{\mu_{Te}}{2U} \right)^{1/2}$$

Classical orbital angular momentum

$$l_c = \sqrt{m_e [(x\dot{y} - y\dot{x})^2 + (x\dot{z} - z\dot{x})^2 + (y\dot{z} - z\dot{y})^2]}$$

Classical magnetic angular momentum

$$m_c = m_e (y\dot{z} - z\dot{y})$$

$$[(n-1)(n-1/2)n]^{1/3} \leq n_c \leq [n(n+1/2)(n+1)]^{1/3}$$

$$l \leq n/n_c l_c < l+1$$

$$(2m-1)/(2l+1) < m_c/l_c < (2m+1)/(2l+1)$$

The total cross sections

$$\sigma = \frac{2\pi b_{max}}{T_N} \sum_j b_j^{(i)}$$

The statistical uncertainty

$$\Delta\sigma = \sigma \left(\frac{T_N - T_N^{(i)}}{T_N T_N^{(i)}} \right)^{1/2}$$

T_N : Total number of trajectories calculated for impact parameters less than b_{max}

$T_N^{(i)}$: Number of trajectories that satisfy the criteria for a given channel

$b_j^{(i)}$: Actual impact parameter for the trajectory corresponding to the channels.

Classical Limits - extension

Improvement of the classical description of the one electron atomic system by including a model potential in the Hamiltonian of the system mimicking quantum features.

Quasi-Classical Trajectory Monte Carlo (QCTMC) Model

$$H_{QCTMC} = T + V_{coul} + V_H$$

Constraining Heisenberg Potential

$$V_H = \sum_{n=a,b} \sum_{i=1}^N f(r_{ni}, p_{ni}; \xi_H, \alpha_H)$$

$$f(r_{\lambda\nu}, p_{\lambda\nu}; \xi, \alpha) = \frac{\xi}{4\alpha r_{\lambda\nu}^2 \mu_{\lambda\nu}} \exp \left\{ \alpha \left[1 - \left(\frac{r_{\lambda\nu} p_{\lambda\nu}}{\xi} \right)^4 \right] \right\}$$

$$\alpha_H - \xi_H$$

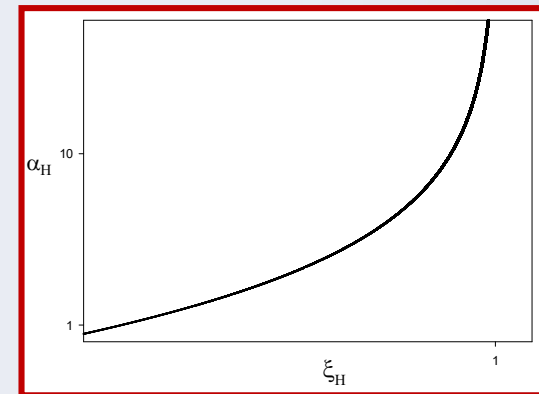
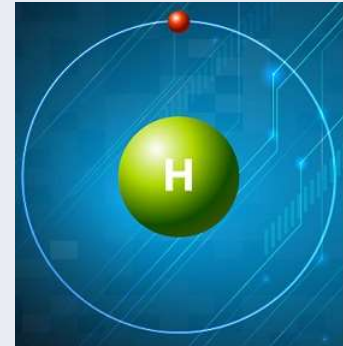
❖ Hamiltonian of hydrogen atom is defined as follows:

$$H = \frac{p^2}{2} - \frac{1}{r} + \left[\frac{\xi_H^2}{4\alpha_H r^2} \right] \exp \left\{ \alpha_H \left[1 - \left(\frac{rp}{\xi_H} \right)^4 \right] \right\}$$

❖ In the ground or lowest-energy configuration, we require $\frac{\partial H}{\partial p} = 0$ and $\frac{\partial H}{\partial r} = 0$

$$E = - \frac{1}{2\xi_H^2 \left(1 + \frac{1}{2\alpha_H} \right)}$$

Electron binding energy = 0.5



Initial conditions for r and p

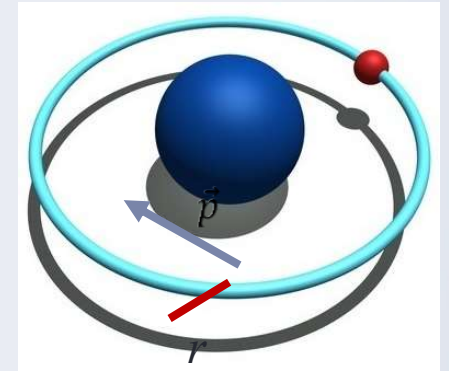
- ❖ In CTMC model, the initial conditions in r and p :

A microcanonical ensemble characterizes the initial state of the target constrained to an initial binding energy of the given shell:

$$\rho_{E_0}(\bar{A}, \dot{\bar{A}}) = K_1 \delta(E_0 - E) = \delta\left(E_0 - \frac{1}{2} \mu_{Te} \dot{\bar{A}}^2 - V(A)\right)$$

$$r_0 = \left| \frac{Z_e Z_T}{2E_b} \right|$$

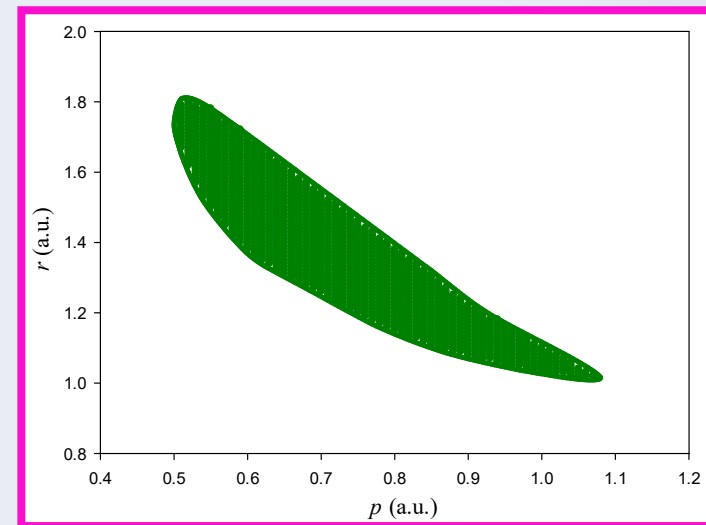
$$p_0 = \sqrt{2|E_b| \mu_{te}}$$



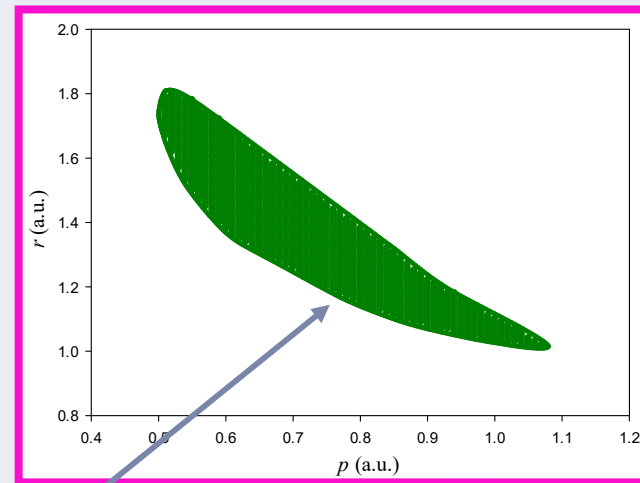
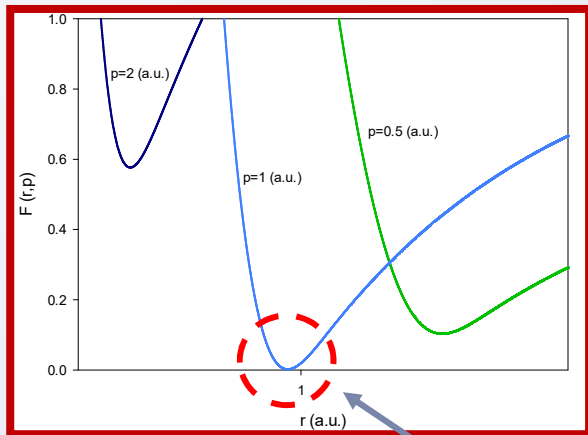
- ❖ In QCTMC model, we considered two conditions that r and p have to satisfy them as follows:

$$\frac{|Z_e Z_T|}{2r} + f_H(r, p) < 0.5$$

$$\frac{p^2}{2\mu_{Te}} - \frac{1}{r} + f_H(r, p) \approx -0.5$$

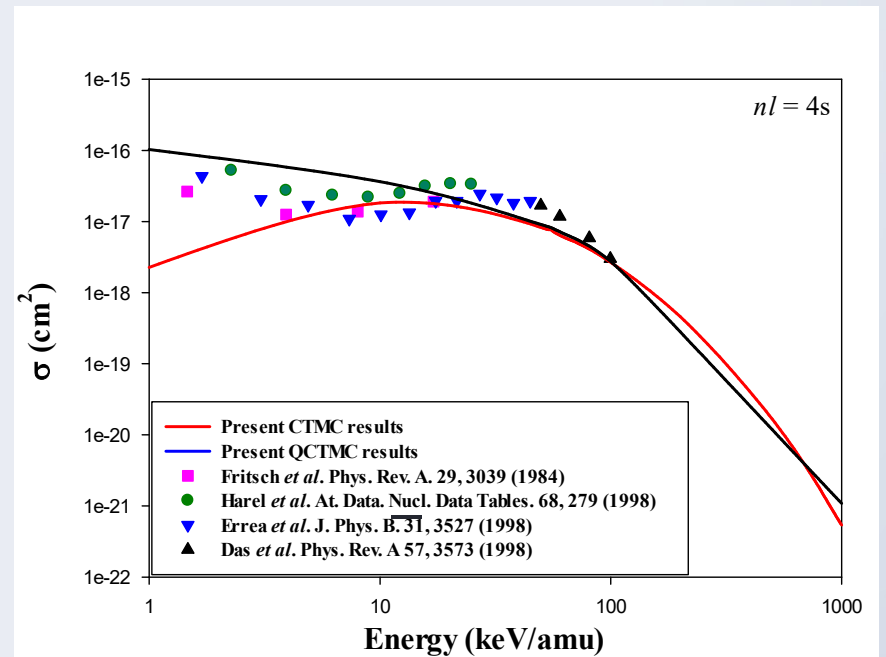
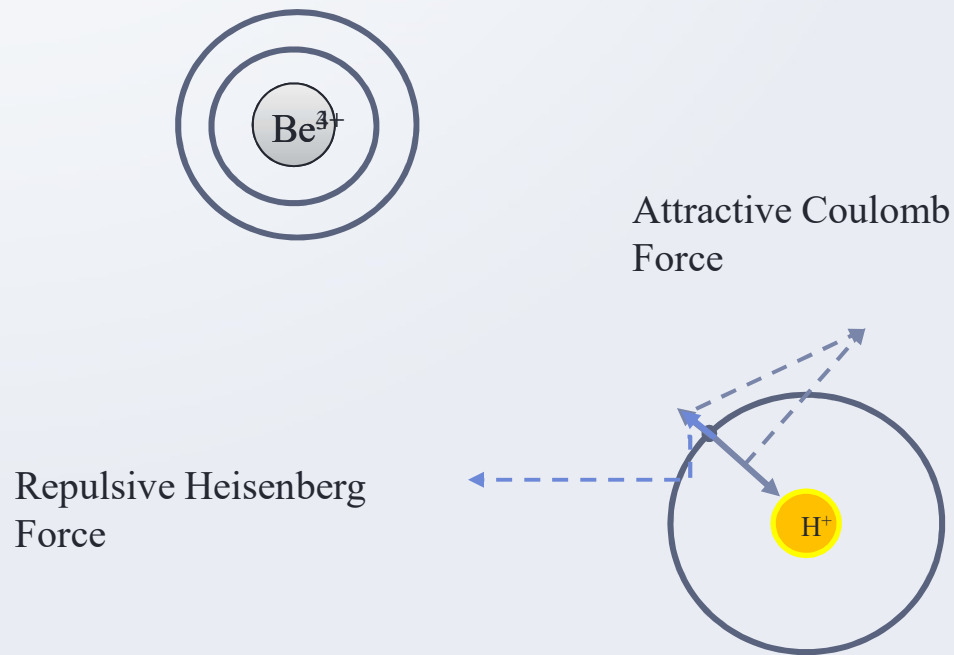


$$F(r, p) = \frac{p^2}{2\mu_{Te}} - \frac{1}{r} + f_H(r, p) + 0.5$$

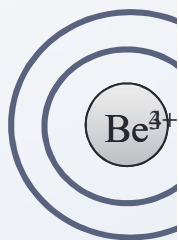


QCTMC initial conditions

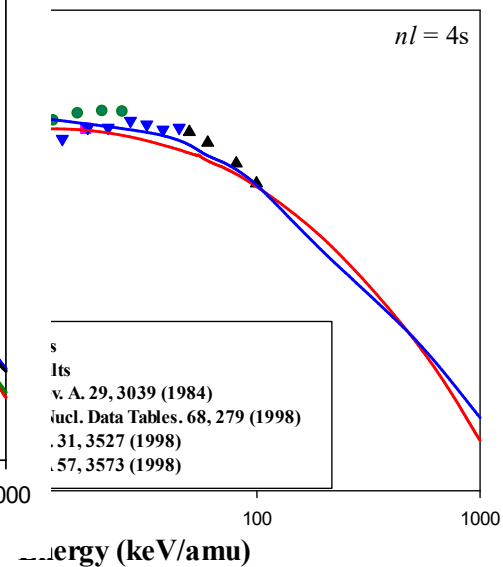
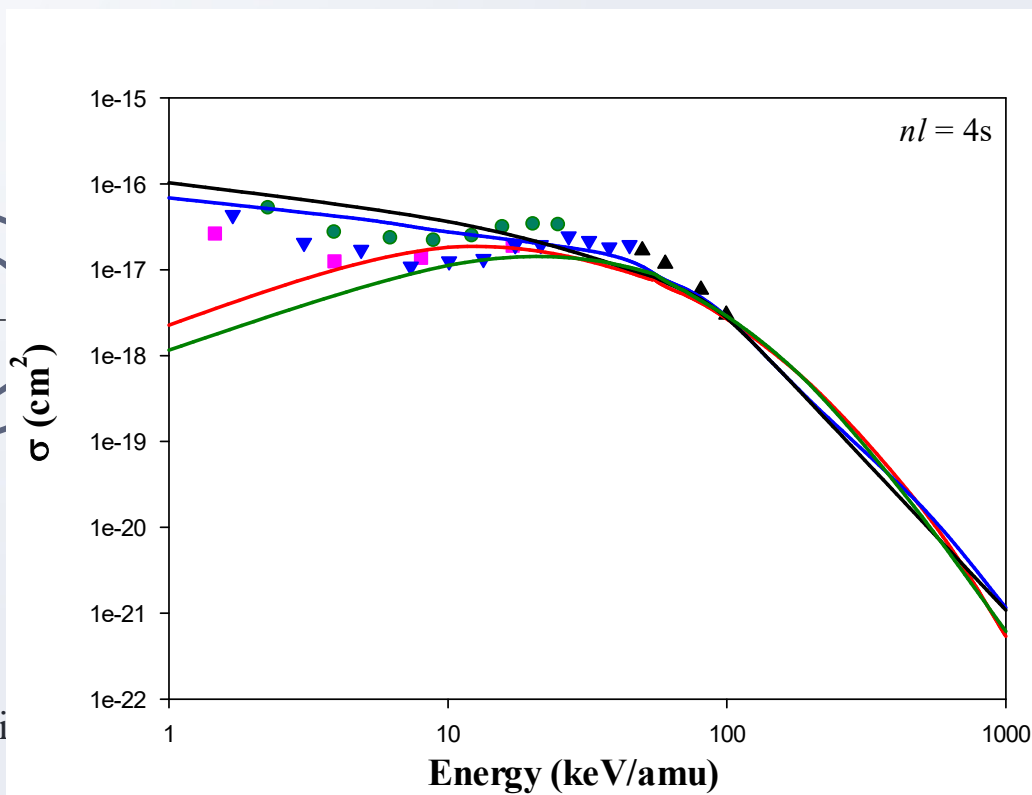
Target-centered scheme



Combined one; i.e., target and projectile-centered scheme



Repulsive Hei
Force

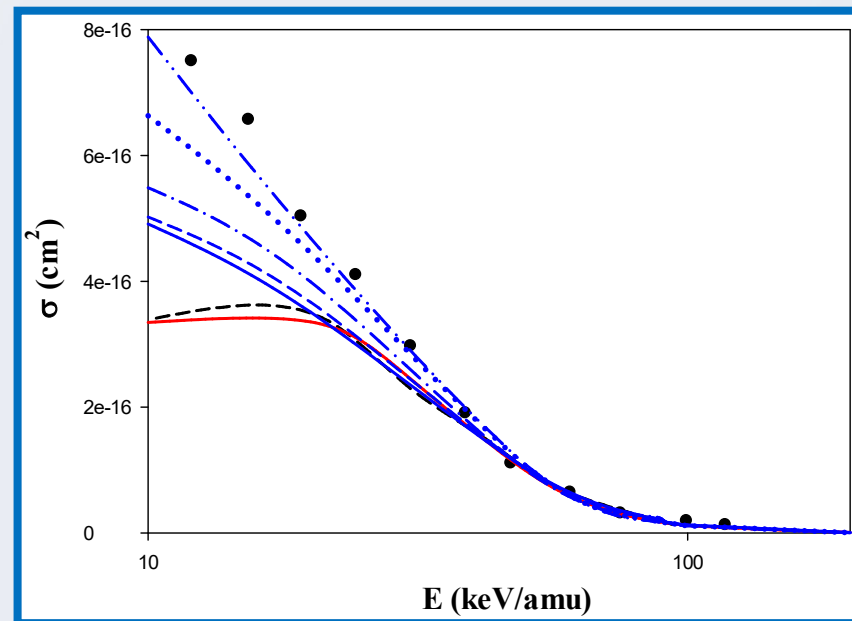


Finding Best Combination of α , ξ

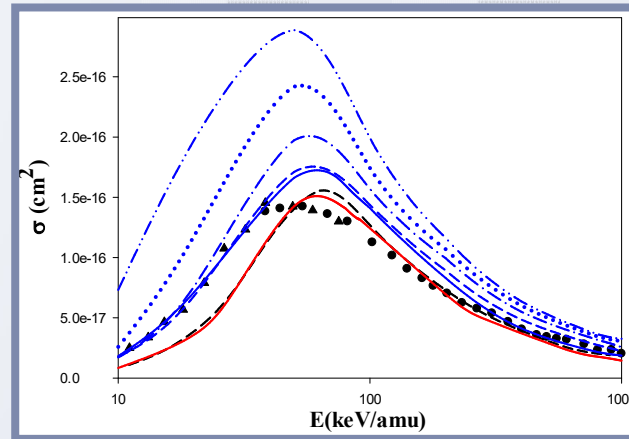
Interaction between H^+ and Hydrogen atom

Total Electron Capture
 $H^+ + H(1s) \rightarrow H + H^+$

- QCTMC ($\alpha_H = 3, \xi_H = 0.9258$)
- - - QCTMC ($\alpha_H = 3.5, \xi_H = 0.9354$)
- · - · QCTMC ($\alpha_H = 4, \xi_H = 0.9428$)
- QCTMC ($\alpha_H = 4.5, \xi_H = 0.9486$)
- · - · QCTMC ($\alpha_H = 5, \xi_H = 0.9534$)
- CTMC
- - - QTMC-EB
- Exp: McClure



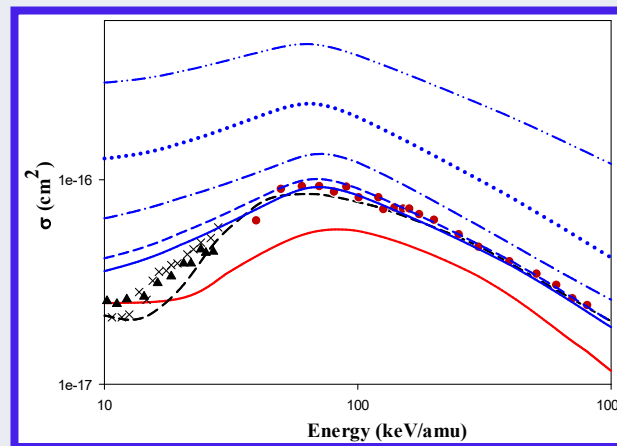
Ionization



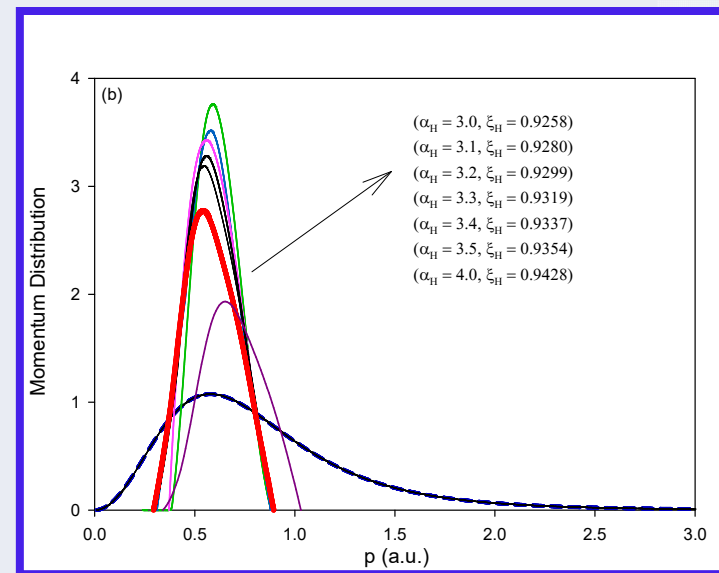
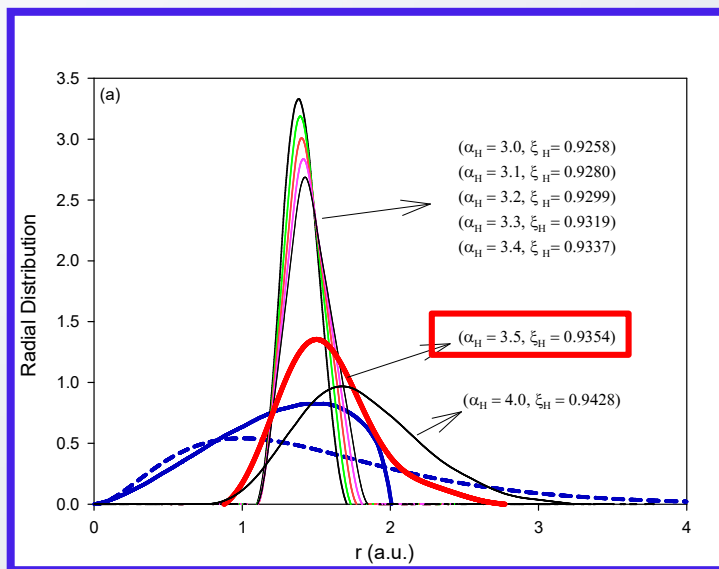
- QCTMC ($\alpha_H = 3, \xi_H = 0.9258$)
- - QCTMC ($\alpha_H = 3.5, \xi_H = 0.9354$)
- . - . QCTMC ($\alpha_H = 4, \xi_H = 0.9428$)
- QCTMC ($\alpha_H = 4.5, \xi_H = 0.9486$)
- . . . QCTMC ($\alpha_H = 5, \xi_H = 0.9534$)
- CTMC
- - - QTMC-EB

- Exp: Shah and Gilbody
- ▲ Exp: Shah and Eliot
- Exp: Detleffsen
- ▲ Exp: Morgan
- ✱ Exp: Kondov

Excitation



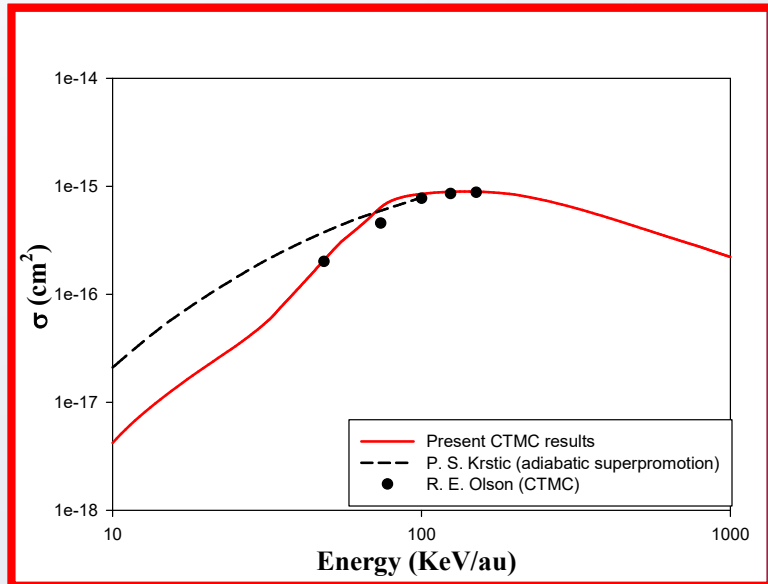
QCTMC Radial & Momentum Distribution



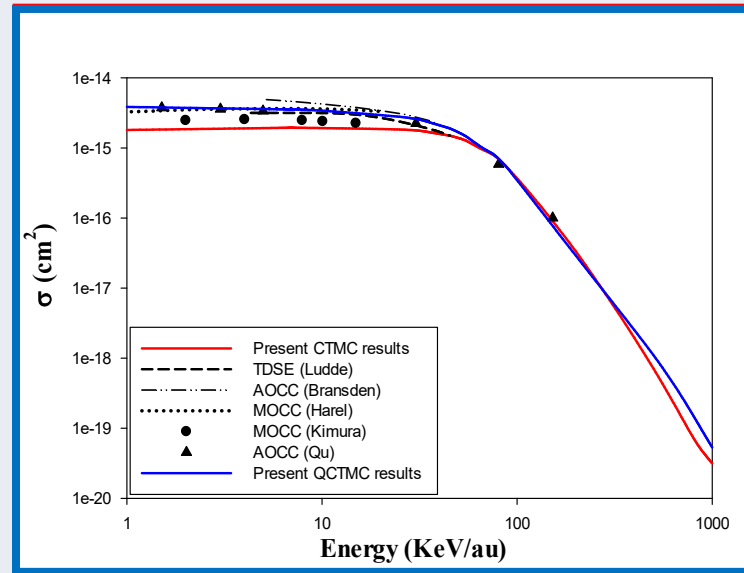
— CTMC Distribution
- - - Quantum Distribution

**Data for Atomic Processes of Neutral
Beams in Fusion Plasma**

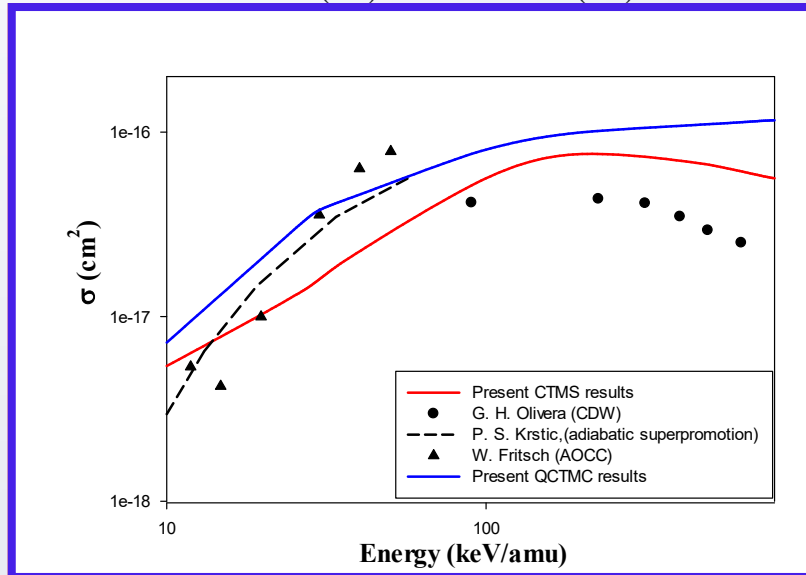
Ionization



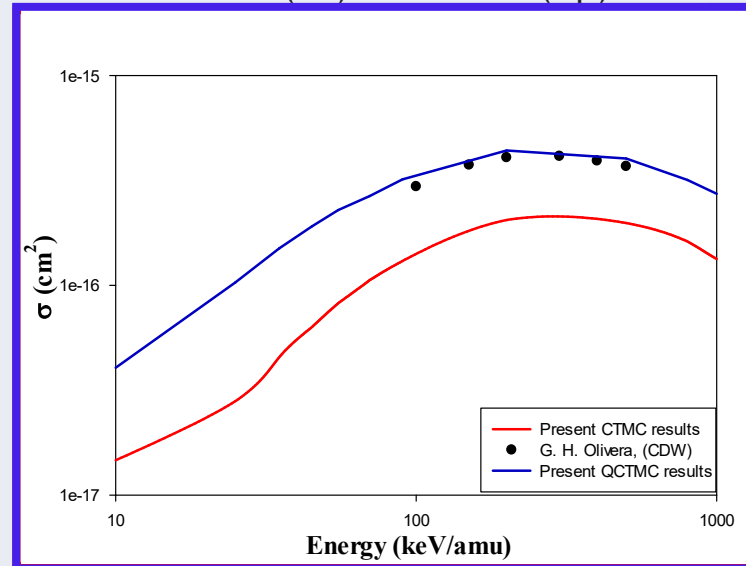
Total Electron Capture



Excitation
 $\text{Be}^{4+} + \text{H}(1s) \rightarrow \text{Be}^{4+} + \text{H}(2s)^*$

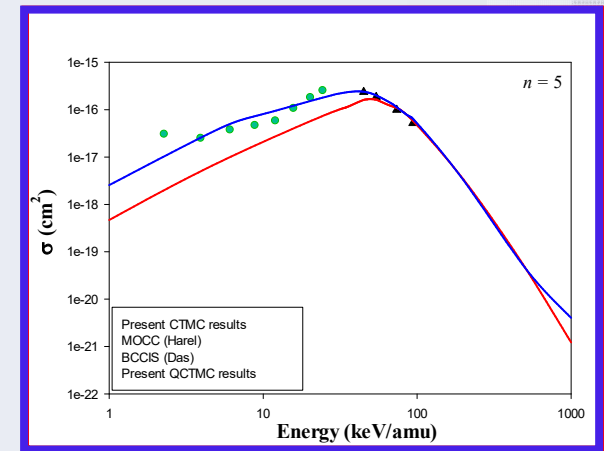
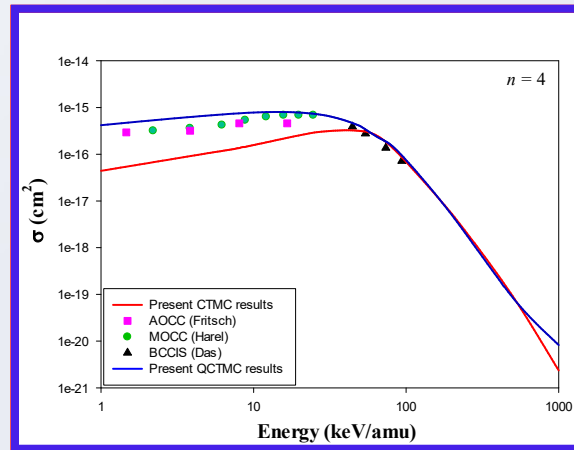
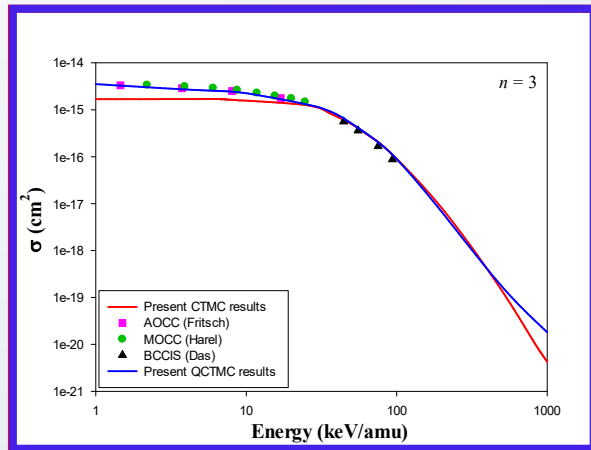


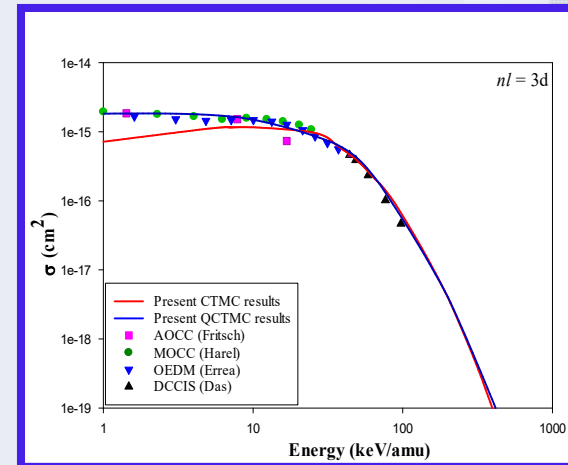
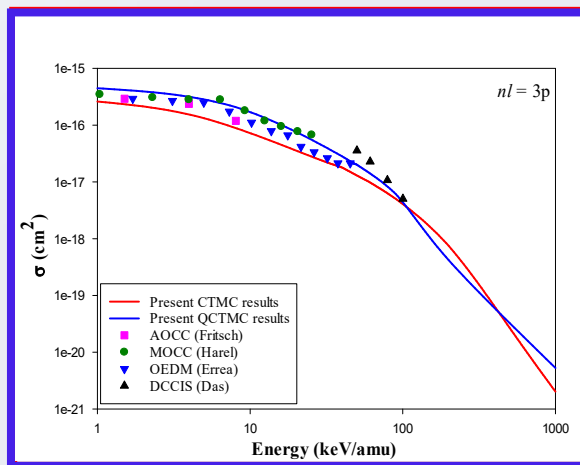
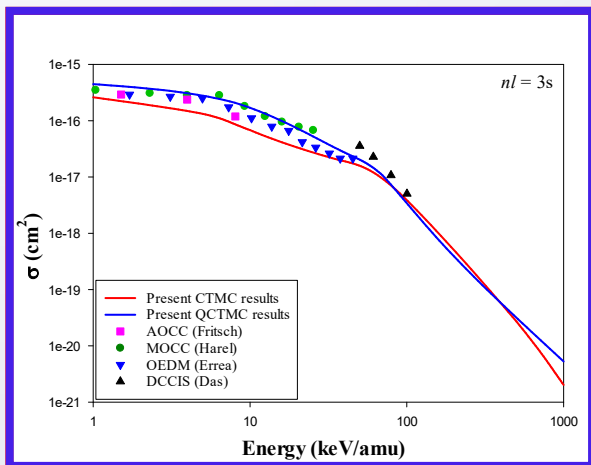
Excitation
 $\text{Be}^{4+} + \text{H}(1s) \rightarrow \text{Be}^{4+} + \text{H}(2p)^*$



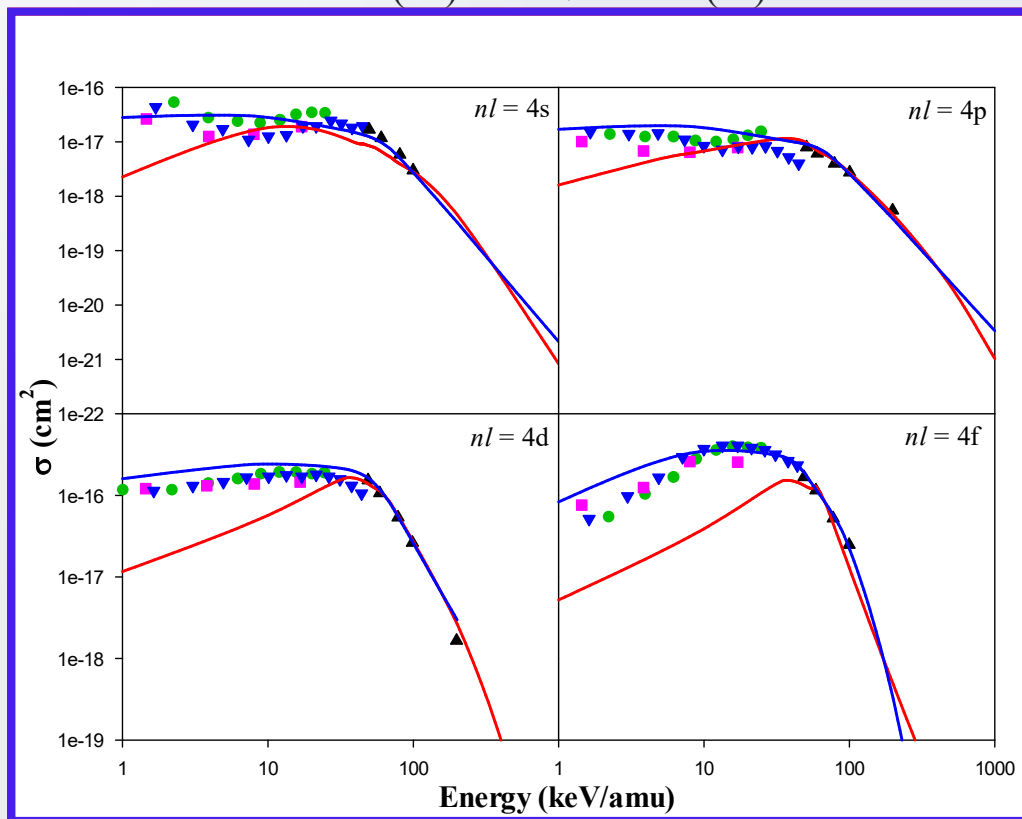
Electron capture cross sections into $n = 3, 4, 5, 6, 8, 10$ and $nl = 3l, 4l, 5l$ states of the projectile in $\text{Be}^{4+} + \text{H}(1s)$ using CTMC and QCTMC models.

State-selective Electron Capture $\text{Be}^{4+} + \text{H}(1s) \longrightarrow \text{Be}^{3+}(n) + \text{H}^+$

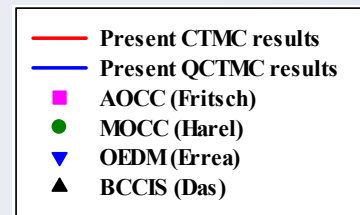
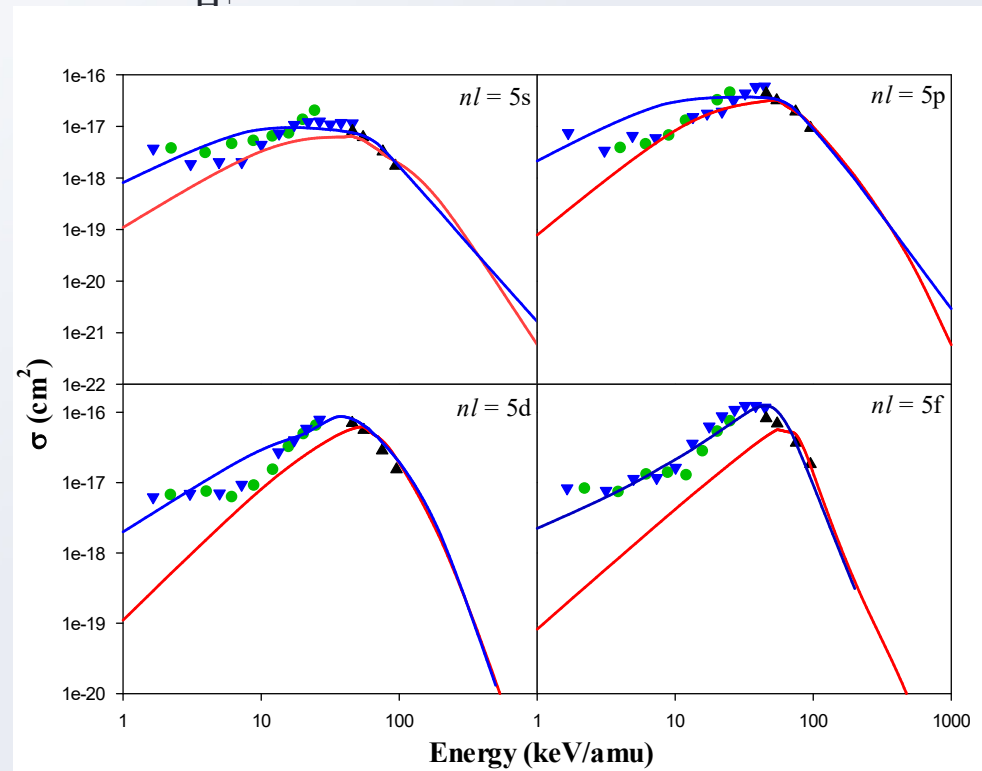


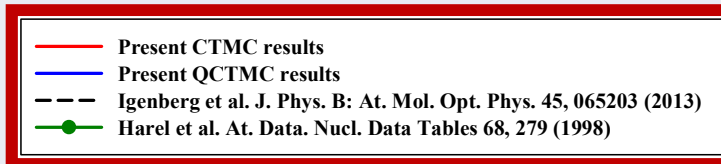
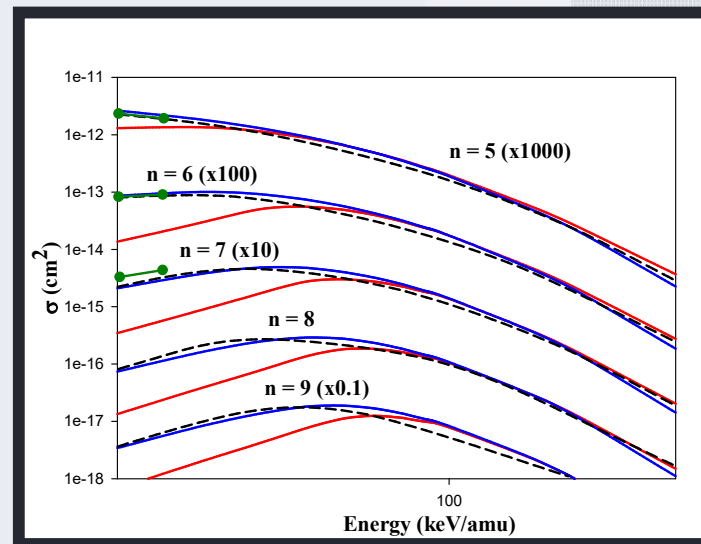
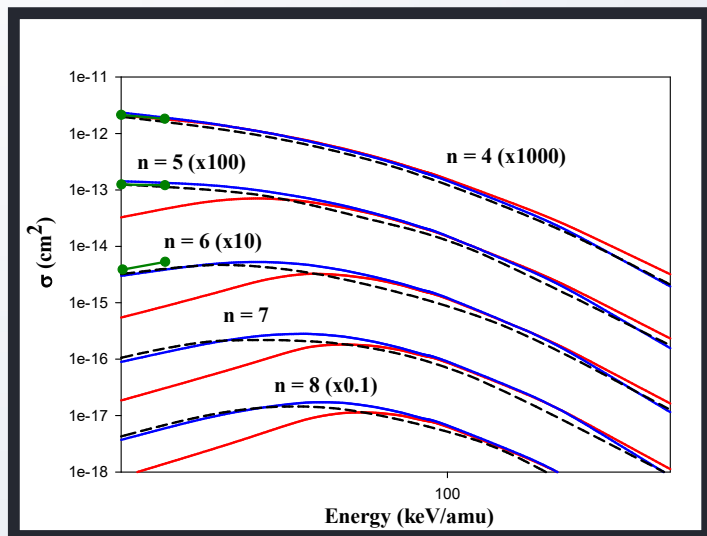


State-selective Electron Capture
 $\text{Be}^{4+} + \text{H}(1s) \longrightarrow \text{Be}^{3+}(4l) + \text{H}^+$



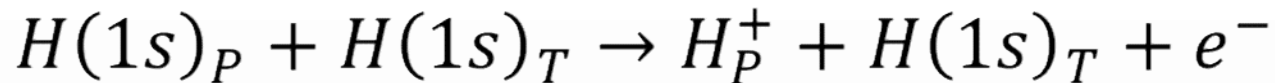
- Present CTMC results
- Present QCTMC results
- AOCC (Fritsch)
- MOCC (Harel)
- ▼ OEDM (Errea)
- ▲ BCCIS (Das)



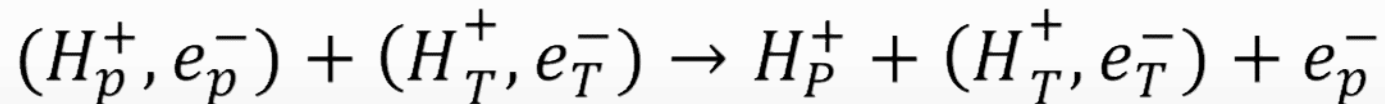


Two ground-state Hydrogen collision System

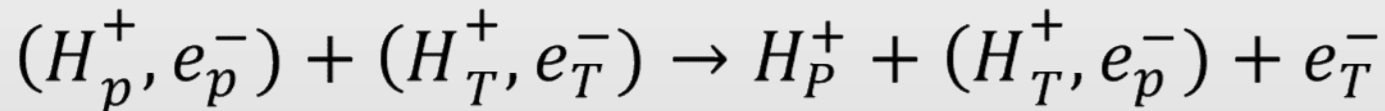
QCTMC Result: Projectile Ionization cross section



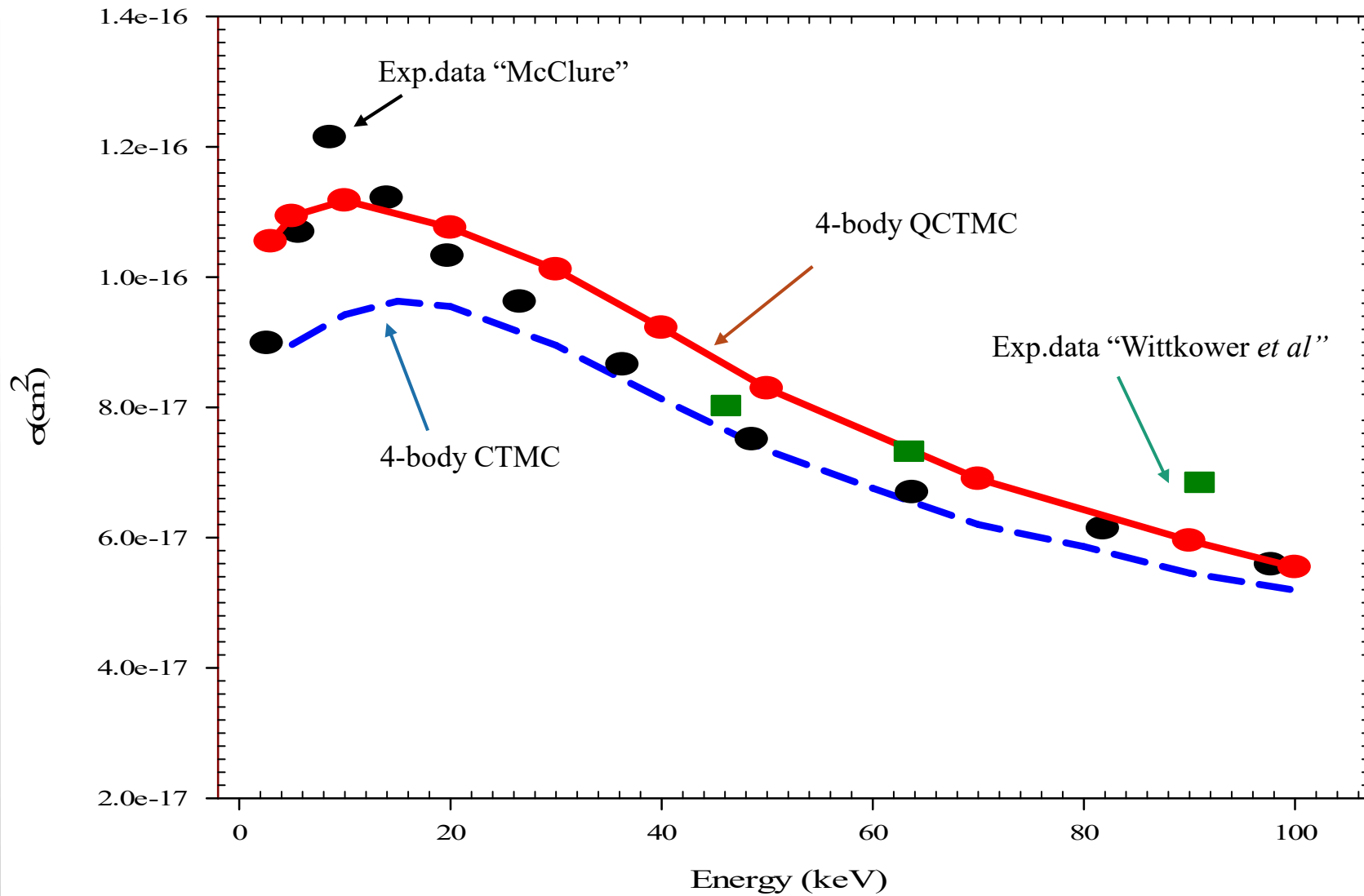
The first possible one the *direct ionization of the projectile* channel. This channel originates from a one-step process



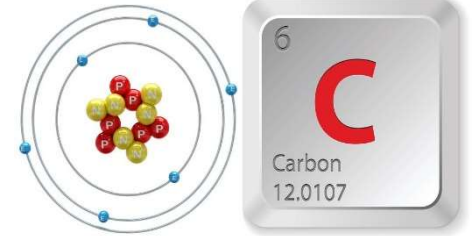
The second possible originates from the multi-electron interaction in a two-step process producing the same final particles.



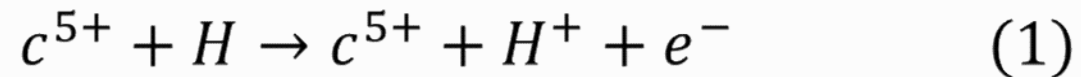
Projectile ionization cross sections



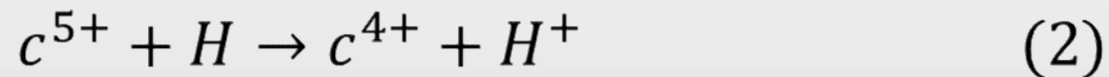
Carbon ions (c^{5+}) with Hydrogen atom collision system



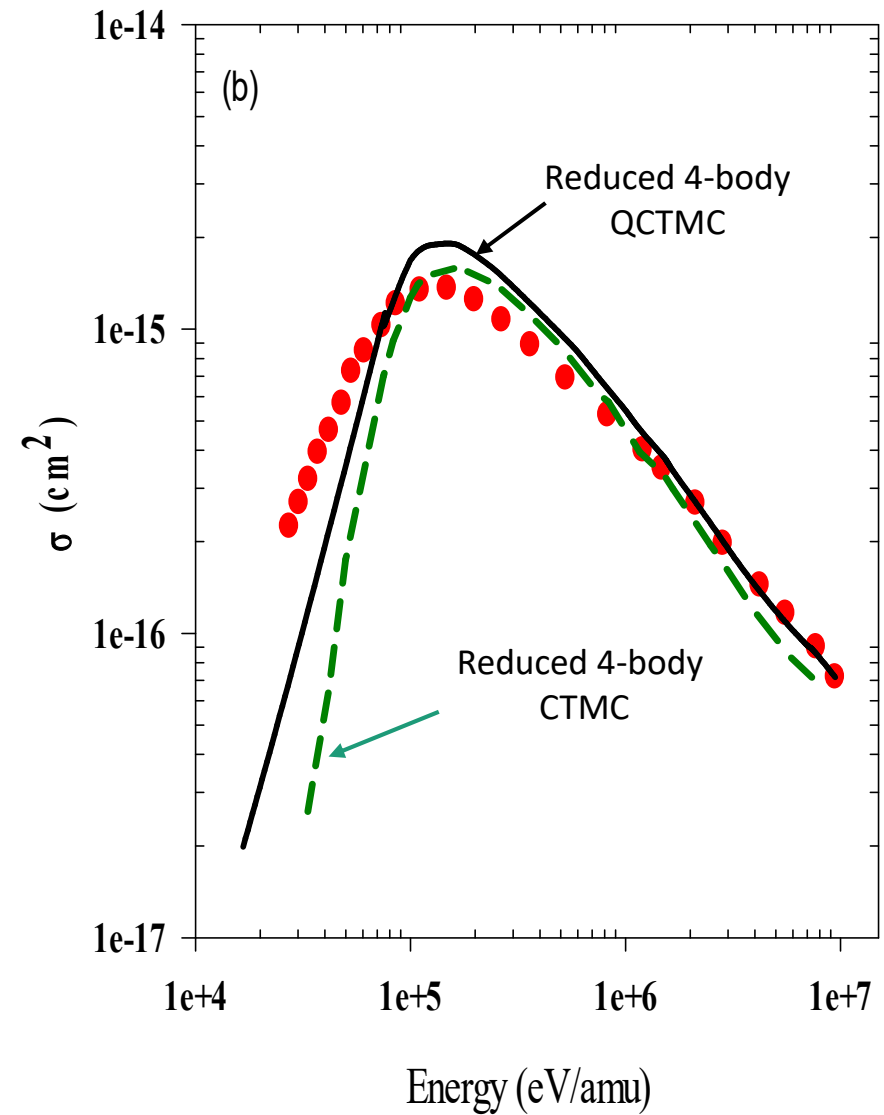
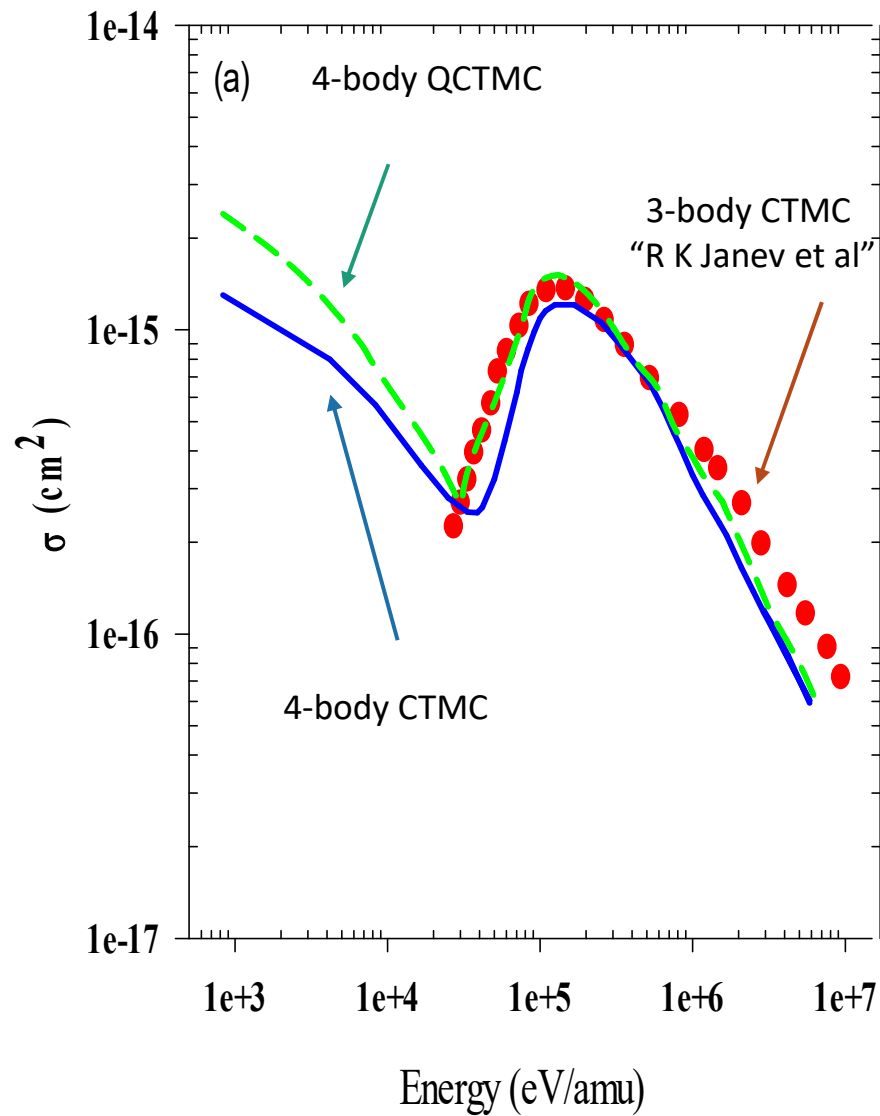
- Target ionization cross section



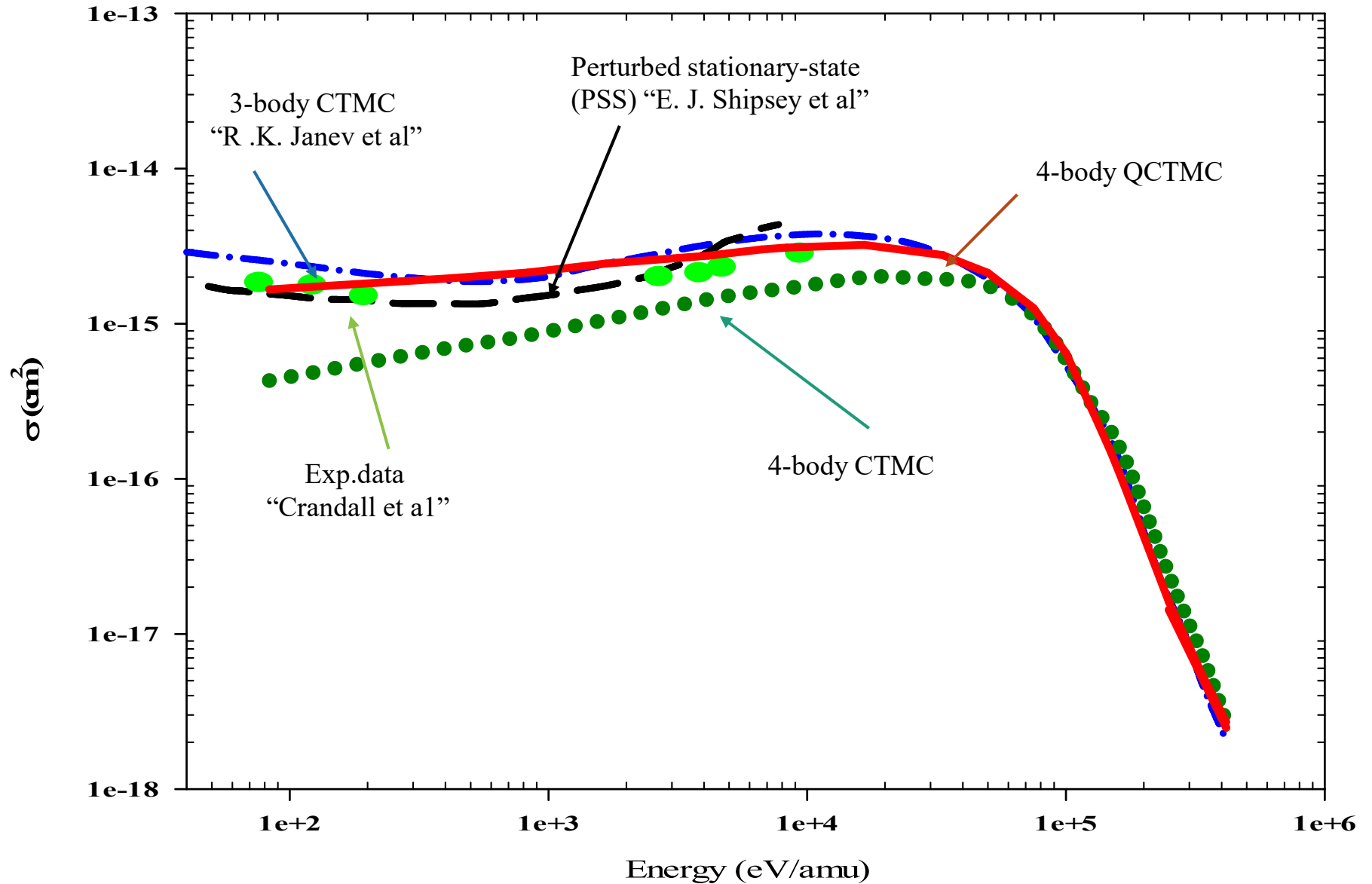
- Electron capture cross section



Target Ionization cross-section



Electron Capture Cross-Section of the Projectile



Publications

Be-H

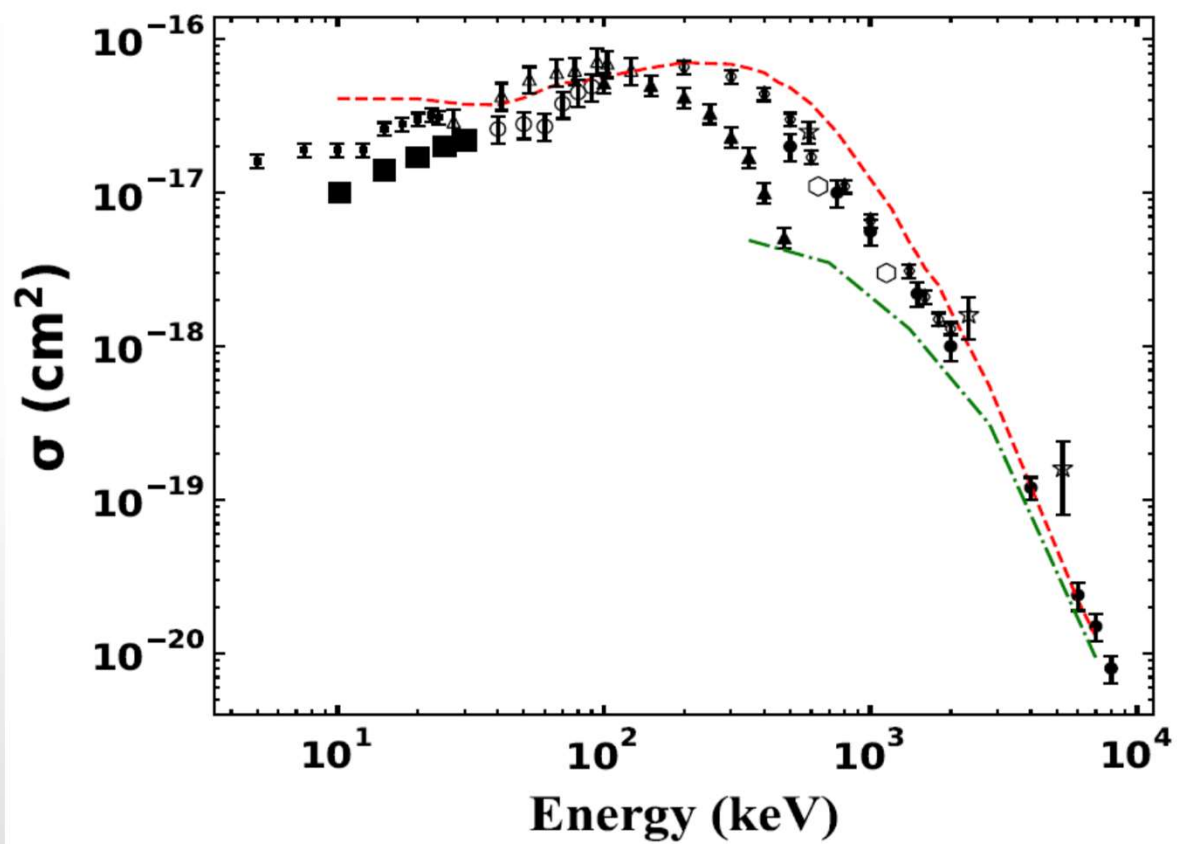
- [1] I. Ziaeeian and K. Tőkési, *Atoms* **8** 27 (2020).
- [2] I. Ziaeeian and K. Tőkési, *EPJD J.* **75** 138 (2021).
- [3] I. Ziaeeian and K. Tőkési, *Sci. Rep.* 20164 (2021).
- [4] I. Ziaeeian and K. Tőkési, *Atoms* **10** 90. (2022)
- [5] I. Ziaeeian and K. Tőkési, , *Journal of Physics B: Atomic, Molecular and Optical Physics* **55** (2022) 245201.
- [6] I. Ziaeeian and K. Tőkési, *Atomic Data and Nuclear Data Tables* **146** 101509 (2022).

H+H type

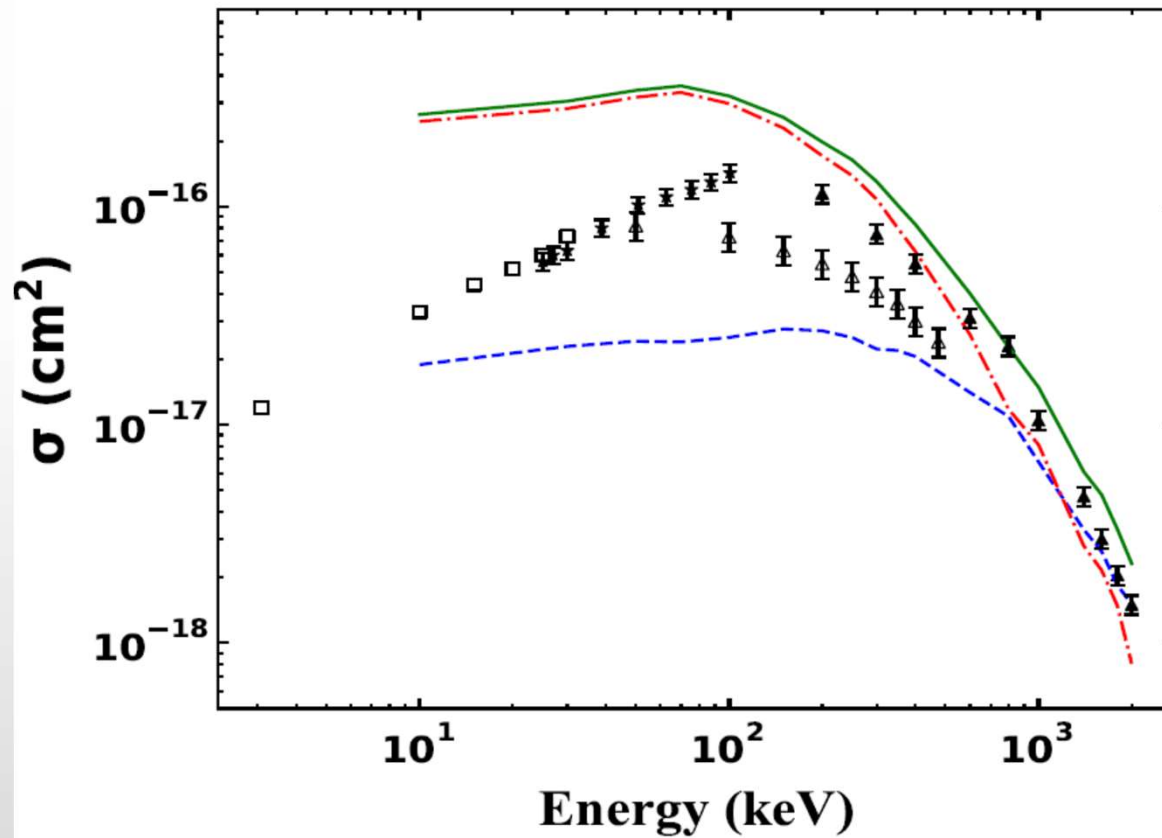
- [1] S.J.A. Atawneh and K. Tőkési, *Atoms* **8** 31 (2020).
- [2] S.J.A. Atawneh and K. Tőkési, *J. Phys. B: At. Mol. Opt. Phys.* **54** 065202 (2021).
- [3] S. J. A. Atawneh and K. Tőkési. *Nucl. Fusion.* **62** 026009 (2021).
- [4] S.J.A. Atawneh and K. Tőkési, *Atomic Data and Nuclear Data Tables* **146** 101513 (2022).
- [5] S. J. A. Atawneh and K. Tőkési. *Phys. Chem. Chem. Phys.* **24** 15280 (2022).

**Atomic Data
for Injected Impurities in Fusion Plasmas**

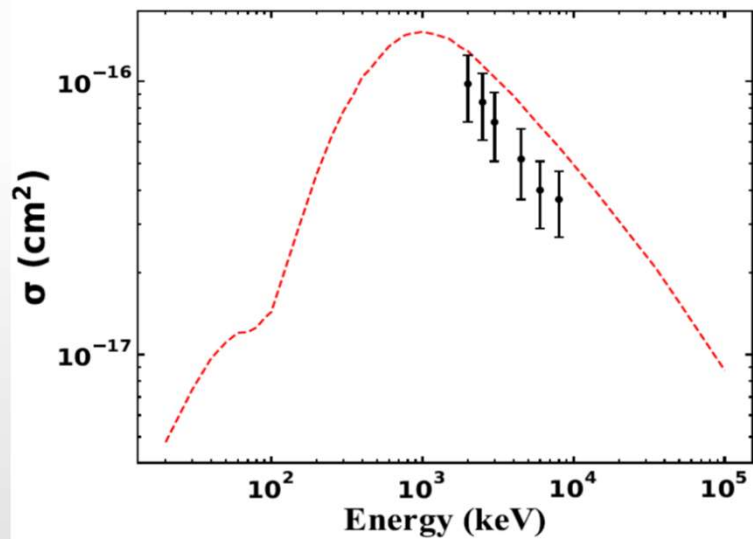
Total cross sections of the single-electron capture from $He(1s)$ by Li^+



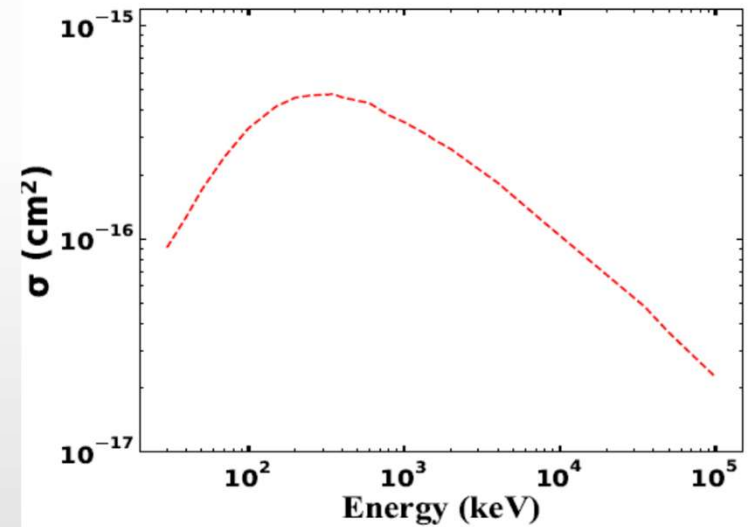
Total cross sections of the single-electron capture from $N(2p)$ by Li^+



**Total cross sections of the
single- electron ionisation from
He(1s) by *Li*⁺**



**Total cross sections of the
single- electron ionisation from
N(2p) by *Li*⁺**



Conclusions

- **Classical method (CTMC) reproduce different experiments for collisions between charged particles and atoms**
- **gives accurate cross sections for ionization, capture, excitation**
- **valid in wide projectile energy range**
- **can describe partial cross sections**
- **QCTMC model, represents one step further towards a better description of the classical atomic collisions. This model with simplicity can time efficiently carry out simulations where maybe the quantum mechanical ones become complicated, therefore, our model should be an alternative way to calculate accurate cross sections and maybe can replace the quantum-mechanical methods.**

Thanks for your
attention!