Plasma interactions with Be surfaces

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US-EU Collaboration on Mixed-Material PMI Effects for ITER US-Japan Technology Exchange Program

PISCES-B experimental arrangement





PISCES-B linear plasma facility at UCSD

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Many pure elemental species plasmas (Ar, Ne, D) and mixed species plasmas (D/He, D/Ar) are possible in the PISCES devices



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Pure D, He, or mixed D/He plasma bombardment produces less Be sputtering than Eckstein's TRIM predicts

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J. Roth et al., FED 37(1997)465.

Fig. 2. Energy dependence of the sputtering yield of Be and BeO bombarded with D at normal incidence. Experimental data [6,13,23–26] and results obtained with computer simulation [6,22].



AES reveals a relatively 'clean' Be surface during sputtering yield measurements

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R.P. Doerner et al. | Journal of Nuclear Materials 257 (1998) 51-58 100 0.1 TRIM-95 (BeO) Be TRIM-95 (Be) (92.1) 0 experiment (plasma spray Be С Sputtering Yield (atoms/ion) experiment (S65C Be) 0 80 0.01 Surface composition (%) Incident ion energy ~100 eV 00 0.001 0 100 200 300 400 500 600 700 800 Sample Temperature (C) 4 1.25.10²² D⁺/cm², -100V bias before dN(E) 20 wattherwith C Fe Cu n after (7.0 (<mark>0.9</mark> O Be 1000 2000 4000 5000 0 3000 500 1000 1500 Time from plasma shut-off (s)



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kinetic energy (eV)

Native beryllium oxide surface is removed early during the plasma exposure





- Distinctive oxygen lines near 777 nm can monitor erosion of O from surface
- Background helium plasma does not change (second order He I line)
- Larger ion energy or flux will remove oxide layer quicker
- BeO (0,0) molecular band emission (@ 470.86 nm) is not detectable



Development of surface morphology decreases erosion



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Initial surface morphology does not significantly affect erosion

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D exposure -50V, < 330K: $j_D = 1 \cdot 10^{22} D_x^+/cm^2$



 \Rightarrow no influence within accuracy of the measurement

UCS

Initial surface morphology does not significantly affect erosion

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D exposure -50V, < 330K: j_D = 1 \cdot 10^{22} D_x^+/cm^2
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 \Rightarrow sub micron morphology develops during exposure UCS

Similar yield evolution with time/fluence is documented in the literature





High flux ion beam Mattox and Sharp, JNM 80(1979)115

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JET-ILW

H. Bergsaker et al., JNM 463(2015)956

Figure 5. SEM images of the surfaces at 1/10 (a) 1/8 (b) and 3/6 (c). All surfaces show elongated structures that are pointing in a direction $15-20^{\circ}$ downwards, with respect to the toroidal direction. The region (c) is slightly depressed with respect to the surrounding layer and has about 4 µm Be deposited, so it seems likely that the cone like structure consists mainly of Be.





D (& He) plasma erosion rates consistently lower than TRIM, Ar plasma sputter data agrees with TRIM





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Distinct surface morphology differences exist after high and low mass sputtering of a Be target

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After Ar plasma sputtering (~75 eV)



After D plasma sputtering (~80 eV)



- Smooth Be surface exists after Ar plasma sputtering
- No significant decrease in Be erosion with fluence during Ar sputtering



Low-Z vs. high-Z sputtering discrepancy is seen from a variety of metals exposed to plasma

- Reason for this behavior is not understood
- Fuel atoms within the stopping distance of surface may influence sputtering/reflection
- Effect is measured for a variety of substrate materials
 - PISCES (Be, Fe, Mo, W, Al)
 - Pilot (Sn)
 - TPE (Be)



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Plasma atoms remaining in the near surface also can reduce the sputtering yield by a factor of 2-4





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Chemically-assisted physical sputtering of BeD is temperature dependent (influence of He is being studied)

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From R. Doerner et al, JNM 390-391 (2009) 681.

Similar e-folding distance of BeD and Be I intensity indicate BeD is physically sputtered, not chemically eroded. Beryllium deuteride is not volatile.

MD simulations of D on Be predicted subsequent erosion measurements

Exp. from D. Nishijima et al, PPCF 50(2008)125007. Sim. from C. Bjorkas et al., New J. Physics (2009).

Temperature Dependent Sputtering : PISCES experiments show qualitative agreement with expectations, but quantitatively show enhanced erosion at high temperature.

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Temperature Dependent Sputtering_: Decrease in average ejection energy of particles accompanies the increase in erosion rate, implying enhanced evaporation. Also true in solids.

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Energetic particle bombardment of surfaces results not only in sputtering, but also in the creation of surface adatoms.

• Adatoms are less strongly bound to the surface and can sublimate at lower temperature, before recombining at surface defect sites [R. P. Doerner et al., J. Appl. Phys. 95 (2004) 4471.]

The total atom loss rate from a surface is written

 $\mathbf{J}_{\text{total}}(\mathbf{T}) = \mathbf{J}_{\text{pl}}\mathbf{Y}_{\text{ps}} + \mathbf{J}_{\text{o}}(\mathbf{T}) + \mathbf{J}_{\text{ad}}(\mathbf{T})$

where, $J_{pl}Y_{ps}$ is the sputtered particle flux

 J_{o} is the flux of lattice atoms sublimated

 J_{ad} is the adatom sublimation flux

And the areal density of adatoms is

$$\mathbf{J}_{\mathrm{o}} = \mathbf{K}_{\mathrm{o}} \mathbf{n}_{\mathrm{o}} \exp^{(-\mathbf{E}_{\mathrm{o}}/\mathrm{T})}$$

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$$\mathbf{J}_{\mathrm{ad}} = \mathbf{n}_{\mathrm{ad}} / \mathbf{t}_{\mathrm{sub}}$$

1

$$n_{ad} = \frac{\mathbf{Y}_{ad} \, \mathbf{J}_{pl}}{\left(\frac{1}{t_{rec}} + \frac{1}{t_{sub}}\right)}$$

Three terms of the total loss rate can be plotted. $J_{total}(T) = J_{pl}Y_{ps} + J_o(T) + J_{ad}(T)$

 Adatom loss term saturates as the fraction sublimated approaches the number created

- At higher temperature the total loss rate is dominated by 'normal' thermal sublimation
- The temperature range where different terms dominate depends on experimental conditions

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Anderl et al.[JNM 273(1999)1] reviewed D retention in Be

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Specimen exposure temperature (°C)

At elevated temperatures, there is little D retention dependence on plasma fluence

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Doerner et al., JNM 257(1998)51 10²² = 100 eV, Flux = 1.8 x 10 21 #/m²-s 200°C : E × 500° C : E_{ion} = 100 eV, Flux = 1.5 x 10²² #/m²-s Deuterium retention (atoms/m²) 10²¹ 10²⁰ 10 0.1 1 Exposure Time (Hrs)

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D retention in clean Be (UHV subML O coverage)

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Reinelt et al. New Journal of Physics 11 (2009) Oberkofler et al. Nuc.l Instrum. Methods 267 (4) (2009)

low flux/fluence implantation at 0.3-3keV with 10¹⁵D cm⁻² s⁻¹:

Only low temperature release peak (~500K) is affected by mixed D/ α He plasma

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Influence of D/He plasma on D retention shows affects only at low exposure temperature

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Retention in Be codeposits will dominate inventory in ITER

[from Roth et al., PPCF 50(2008)103001]

Retention in codeposited Be can be predicted based on codeposition conditions

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D/Be = $(5.82 \times 10^{-5})E_n^{1.17\pm0.2} (\Gamma_D/\Gamma_{Be})^{-0.21\pm0.1} \exp^{(2273\pm311/T_c)}$

[from G. De Temmerman et al., NF 48(2008)075008]

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Release of D from Be codeposits : passivating oxide layer at interface has little impact on release behavior

- Double layer codeposit release agrees with sum of two ٠ single layers (thickness ~ 100 nm) collected on PISCES-B witness plates
- Slight increase in release temperature with increasing • codeposit thickness
- Difficult to model due to small differences in each codeposit

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D release from single and double layer, water cooled codeposits

A single TMAP model describing Be-D co-deposit release agrees with experiments over a wide range of parameter space

- Identical Be-D co-deposits on W spheres is ideal for modeling D release
 - Reproduces shift in release to higher temperature with increasing layer thickness (TMAP red)
 - Reproduces change in release during differing heating rate bakes (TMAP green)
 - Reproduces release decrease during temperature holds
 - Reproduces release behavior during subsequent temperature increase after hold (TMAP blue)

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ITER will likely encounter depleted codeposit under layers covered by saturated codeposit over layers

Multi-layered co-deposit in TEXTOR from M. Rubel et al., Phys. Scr. T103(2003)20.

- Periodic baking of ITER will be used to control in-vessel tritium accumulation (depleting codeposits) [definition – under layers]
- Subsequent operation will lay down new, saturated Be-D/T codeposits [definition – over layers]
- Under layers could be fully, or partially, depleted of D/T
- How does depleted under layer affect D/T release from over layer, or does T diffuse and retrap in previously depleted trap sites?

Agreement between the standard TMAP model and experiment is no longer achieved

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Formation and growth of an oxide layer during baking has a profound effect on D release behavior

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Preferential loss of Be from the surface of mixed Be-W targets (from Romania) occurs during D and D/He plasma exposure

- Samples prepared by thermoionic arc (TVA) process in Romania
- Ion energy is above Be sputtering threshold, but below W sputtering threshold
- Be I line emission quickly decreases from exposed targets
- W-rich surfaces result after D or D/He plasma exposure regardless of initial surface composition
- Similar to Fe/W behavior in RAFM steels

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Loss of surface Be results in W fuzz growth at high temperature

[from I. Jepu et al., submitted to JNM 2016]

D retention in Be-rich TVA layers follows that of bulk Be, whereas W-rich TVA layers follows that of bulk W

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D retention • measured by D₂ plasma TDS after 473 D₂-0.1He plasma 5x10²¹ K exposure No measureable • Retention in D desorbed _D (m^{-2}) bulk W after D, after 1073 K or D/0.1He, -0exposure plasma exposure [Baldwin NF Retention in bulk Be 51(2011)103021] [A] after D, or D/0.1He, X 0 [B] plasma exposure [Baldwin PSI 2016] 0.0 0.5 1.0

Nominal W fraction

[from I. Jepu et al., submitted to JNM 2016]

Understanding pre-created Be/W surface layers is important, but plasma-created Be/W layers will also form when the incident plasma on W contains Be.

In PISCES-B, this can be investigated by using the Be oven to seed controlled amounts of be impurities into the plasma column.

Two cases are possible, when the incident flux of Be is less than the erosion rate of Be from the surface, or when the incident flux is greater than erosion.

What will be the result of Be layers on plasma exposed W surfaces?

Weight Percent Tungsten 0 40 60 70 80 90 95 100 3500 3422°C 3000 2500 Temperature ^oC 2100±50°C 1289°C 2000 ~95 -60 <1750°C 1500 (BBe) 1000 (αBe) 500 $\dot{20}$ 30 <u>90</u> 40 50 80 60 70 100 Be Atomic Percent Tungsten W

From H. Okamoto and L.E. Tanner, in "Phase Diagrams of Binary Tungsten Alloys", Ed. S.V. Naidu and P. Rao, Indian Institute of Metals, Calcutta, 1991.

 Resulting alloys (Be₂₂W and Be₁₂W) have much lower melting temperatures (~1500°C)

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 Be₂W alloy also has somewhat lower melting temperature (<2250°C)

Be-containing plasma on W forms a thin Be_2W surface layer but does not drastically increase retention in W

A thin (few nm) Be₂W layer forms on the surface of W exposed to D+Be plasma (when the Be flux is smaller than erosion), which does not change retention significantly, but may act to inhibit release of implanted D

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Laser transients on Be-coated W create a rich variety of morphology changes

Macroscopic Be layers grow on W targets when the Be arrival rate is greater than the erosion rate

Ablation of Be, recrystallization of W
Be-W alloying and cracking
Be melt flow
Delamination and holes in Be coating

Sputter XPS depth profiles show Be-W mixture at center of spot

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Sacrificial Be layer

Can eroded wall material that is deposited on the divertor provide protection to the underlying W? How thick does Be layer need to be?

Energy density required to sublimate (or equivalently, melt and evaporate) a layer of Be with thickness d: $\mathcal{E}_{Be} = \Delta H \rho d / \mu$ "Be formation energy density"

 ρ is density, μ is molar mass, and enthalpy of formation for Be is Δ H = 324 kJ/mol (*CRC Handbook*)

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SEM images support simple estimate for Be layer removal

laser

Plasma response to laser heating of Be targets differs from that of W, Mo targets

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- Be, W and Mo targets were irradiated with 20 msec long laser pulses (max. energy 50 J) and the response of the plasma was observed
- Penetration distance (or decay length) of ablated material was measured along with background plasma line emission
- PIC simulations show Mo and W atoms are localized close to the targets due to their slow ejection velocity and higher ionization cross-section, whereas ablated Be penetrates more deeply into the plasma column causing cooling of the plasma

[from K. Ibano et al., PSI 2016]

Overview of experiments at PISCES-B

Samples were irradiated with shots of the maximum laser energy.

 \rightarrow Once a shielding behavior was observed, smaller energy shots were tested.

Analysis of experimental results

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$$f_{(x)} = I_{\text{peak}} * e^{-\lambda x_{\text{peak}}}$$

Longer Decay length

- → Longer mean free path
- \rightarrow Faster particle ejection

and/or

 \rightarrow Lower plasma temperature

Experimental observations: W and Mo

Experimental observations: Beryllium

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Longer decay length during laser shots means longer mean free path of Be vapor, but lower vapor velocity implies lower plasma temperature.

(Vapor Shielding)

Experimental observations: Beryllium

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Time [ms]

In addition to ablation velocity, atomic physics differs between W and Be

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Ionization cross section Radiated power 10⁻¹² 10⁻³⁰ Be⁺⁰ n_e=10¹⁹ m⁻³ 10¹⁹ m⁻³ Be⁺¹ 10⁻³¹ Be⁺² Be⁺³ 10⁻¹³ radiation power (J-m³/s) W⁺⁰ 10⁻³² W⁺¹ sigmav (m³/s) W⁺² W⁺³ 10⁻¹⁴ 0-33 Be⁺⁰ Be⁺ 0-34 Mo^{+0} 10⁻¹⁵ Mo⁺ 10⁻³⁵ W⁺⁰ **\//**+` 10⁻¹⁶ 10⁻³⁶ 100 1000 10000 10 10 100 Temperature (eV) 0.1 1 1000 10000 Temperature (eV)

PIC simulations predict more energy dissipation in PISCES by Be than W

- Be impurities significantly reduce the heat flux to the wall.
- In comparison, no apparent reduction was observed for W.

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- High-flux of low-z ions still appears to reduce erosion of surfaces (JET-ILW sees lower Be erosion than PISCES-B) chemistry and temperature dependent losses may be important
- Fuel retention will be dominated by codeposition with Be and estimates seem robust, although impurities could alter values
- Release of D from Be is well modeled by TMAP (not including impurity effects)
- Be/W mixed material formation on hot divertor plates is still an open issue during both steady-state as well as transient events
- The response of the plasma to particle ejection from various materials has shown differences between Be and W

While much progress has been made in PMI with Be, there are still many details to be investigated

