

Consiglio Nazionale delle Ricerche

Ab-initio electron-molecule state-resolved dynamics for non-equilibrium plasmas modelling

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25th Data Center Network meeting, IAEA, Vienna, Austria – 1 October 2019

Plan of the presentation:

- 1. Theory: electron-molecule scattering
- *ab-initio* state-resolved cross sections:
 - A. Nitrogen and Oxygen: Atmospheric (re-)entry and hypersonic
 - **B.** Carbon dioxide: Plasma-based CO₂ conversion
 - **C. BeH**⁺, **H**₂⁺, **OH**, **ArH**⁺: I.T.E.R and controlled fusion; Astrochemistry
- 3. Conclusions and perspectives

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The R-matrix method





J. Tennyson, Phys. Rep., **491**, 29 (2010)

The R-matrix method

Inner region: N+1 electrons; *ab-initio* quantum chemistry codes, including exchange correlations; continuum orbitals

$$u_i(a) = \sum_j R_{ij}(a, E) [r \, u_j' - b \, u_j]_{r=a}$$







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$$u_i(a) = \sum_j R_{ij}(a, E) [r \, u_j' - b \, u_j]_{r=a}$$

$$R_l(a, E) = \frac{1}{2} \sum_{\lambda} \frac{w_{\lambda l}(a)^2}{E_{\lambda} - E}$$

Inner region

Outer region: exchange and electron-electron correlations are neglected and it need to take into account only long–range multi–polar interactions between the scattering electron and the target



$$u_l(r) \sim \frac{1}{\sqrt{k}} (\sin \varphi_l + \cos \varphi_l K_l)$$

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$$u_l(r) \sim \frac{1}{\sqrt{k}} (\sin \varphi_l + \cos \varphi_l K_l)$$

Basic formulas for electron-molecule scattering

Processes occurring for resonant electron-molecule scattering:

$$e^{-} + AB(v) \rightarrow AB^{-} \rightarrow AB(w) + e^{-}$$
$$\rightarrow A^{-} + B$$
$$\rightarrow A + B + e^{-}$$
$$e^{-} + AB^{+}(v) \rightarrow AB^{*} \rightarrow A + B$$

Vibrational Excitation Dissociative Attachment Dissociative Excitation Dissociative Recombination



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Vibrational Excitation Dissociative Attachment Dissociative Excitation Dissociative Recombination

e

AB

$$\left(T_N + V^- - \frac{i}{2}\Gamma - E\right)\xi(R) = -V_{dk}^*(\epsilon, R)\,\chi_i(R)\,,$$
Neutral molecule

Resonant

Vibrational-excitation cross section

 $e + AB(v) \rightarrow AB^{-} \rightarrow AB(v') + e$



J. N. Bardsley and F. Mandl, Reports on Progress in Physics, **31**, 471 (1968)



Resonant Vibrational-excitation cross section

 $e + AB(v) \rightarrow AB^{-} \rightarrow AB(v') + e$

J. N. Bardsley and F. Mandl, Reports on Progress in Physics, **31**, 471 (1968)



J. N. Bardsley and F. Mandl, Reports on Progress in Physics, **31**, 471 (1968)



Dissociative recombination



Dissociative recombination





Dissociative recombination

A Giusti, Journal of Physics B: Atomic and Molecular Physics, **13**, 3867 (1980)

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electron-N₂ resonant scattering

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The resonance at 2.3 eV in electron-N₂ scattering is described in term of the resonant state $N_2^-(X \, {}^2\Pi_g)$

The quantum chemistry codes MOLPRO and UK-R-Matrix have been used to calculate potential energy curves, resonance width and scattering

cc-pvQZ basis set MR-CI model

N₂ target: $(1\sigma_g, 1\sigma_u)^4 (2\sigma_g, 2\sigma_u, 1\pi_u, 3\sigma_g, 1\pi_g, 3\sigma_u)^{10}$

e+N₂ scattering: $(1\sigma_g, 1\sigma_u)^4 (2\sigma_g, 2\sigma_u, 1\pi_u, 3\sigma_g, 1\pi_g, 3\sigma_u)^{11}$ $(1\sigma_g, 1\sigma_u)^4 (2\sigma_g, 2\sigma_u, 1\pi_u, 3\sigma_g, 1\pi_g, 3\sigma_u)^{10} (2\pi_u, 4\sigma_g, 5\sigma_g, 2\pi_g, 1\delta_g, 4\sigma_u)^1$

> V. Laporta *et al.*, Plasma Sources Sci. Technol. **23**, 065002 (2014) V. Laporta *et al.*, Plasma Sources Sci. Technol. **21**, 055018 (2012)



V. Laporta *et al.*, Plasma Sources Sci. Technol. **23**, 065002 (2014) V. Laporta *et al.*, Plasma Sources Sci. Technol. **21**, 055018 (2012)

Comparison with experiments



 $e^- + \mathrm{N}_2(\mathrm{X}\,{}^1\Sigma_g^+;\,\nu,J) \longrightarrow \mathrm{N}_2^-(\mathrm{X}\,{}^2\Pi_g) \longrightarrow e^- + \mathrm{N}_2(\mathrm{X}\,{}^1\Sigma_g^+;\,\nu',J)$

V. Laporta *et al.*, Plasma Sources Sci. Technol. **23**, 065002 (2014) V. Laporta *et al.*, Plasma Sources Sci. Technol. **21**, 055018 (2012)

Vibrational-excitation process:

Comparison with experiments



 $e^- + \mathrm{N}_2(\mathrm{X}\,{}^1\Sigma_g^+;\,\nu,J) \longrightarrow \mathrm{N}_2^-(\mathrm{X}\,{}^2\Pi_g) \longrightarrow e^- + \mathrm{N}_2(\mathrm{X}\,{}^1\Sigma_g^+;\,\nu',J)$



Vibrational-excitation process:

Electron-impact dissociation:

$$e^- + \mathrm{N}_2(\mathrm{X}\,{}^1\Sigma_g^+; \nu, J) \longrightarrow \mathrm{N}_2^-(\mathrm{X}\,{}^2\Pi_g) \longrightarrow e^- + 2 \mathrm{N}({}^4\mathrm{S})$$



electron- $O_2(X, V)$ resonant scattering

electron- $O_2(X, v)$ resonant scattering

In order to describe the low-energy electron- $O_2(X^3\Sigma_g^-)$ resonant scattering it needs to include in the calculations four resonant states, ${}^2\Pi_g$, ${}^2\Pi_u$, ${}^4\Sigma_u^-$, ${}^2\Sigma_u^-$ of O_2^-

Potential energy curves and resonance widths were obtained by MOLPRO and R-matrix within **aug-cc-pvQZ** basis-set and **MR-CI** model

Orbital configurations for **O**₂ target:

$$\begin{array}{ccc} (2a_g, 1b_{1u})^6 & (3a_g, 2b_{3u}, 2b_{2u}, 3b_{1u}, 1b_{2g}, 1b_{3g})^{10} \\ 3 \ \text{core orbitals} & 9 \ \text{valence orbitals} \\ \text{frozen electrons} \end{array}$$

Two orbital configurations for $e + O_2$ scattering: $(2a_g, 1b_{1u})^6 (5a_g, 2b_{3u}, 2b_{2u}, 4b_{1u}, 2b_{2g}, 2b_{3g})^{11}$ and $(2a_g, 1b_{1u})^6 (5a_g, 2b_{3u}, 2b_{2u}, 4b_{1u}, 2b_{2g}, 2b_{3g})^{10} (6a_g, 3b_{3u}, 3b_{2u}, 1b_{1g}, 5b_{1u}, 3b_{2g}, 3b_{3g}, 1a_u)^1$



V. Laporta, R. Celiberto and J. Tennyson, *Phys. Rev.* A91, 012701 (2015). V. Laporta, R. Celiberto and J. Tennyson, *Plasma Sources Sci. Technol.* 22, 025001 (2013)

Vibrational-Excitation process:

$$e^- + O_2(X^3\Sigma_g^-; \nu) \rightarrow O_2^-({}^2\Pi_g, {}^2\Pi_u, {}^4\Sigma_u^-, {}^2\Sigma_u^-) \rightarrow e^- + O_2(X^3\Sigma_g^-; \nu')$$



Vibrational-Excitation process:

$$e^- + O_2(X^3\Sigma_g^-; v) \rightarrow O_2^-({}^2\Pi_g, {}^2\Pi_u, {}^4\Sigma_u^-, {}^2\Sigma_u^-) \rightarrow e^- + O_2(X^3\Sigma_g^-; v')$$

Full set of *Vibrational Excitation* cross sections for *j* = 1...



Dissociative-Electron-Attachment

$$e^{-} + O_2(X^{3}\Sigma_{g}^{-}; v) \rightarrow O_2^{-}({}^{2}\Pi_{g}, {}^{2}\Pi_{u}, {}^{4}\Sigma_{u}^{-}, {}^{2}\Sigma_{u}^{-}) \rightarrow O({}^{3}P) + O^{-}({}^{2}P)$$



Electron energy (eV)

V. Laporta, R. Celiberto and J. Tennyson, *Phys. Rev.* A91, 012701 (2015). V. Laporta, R. Celiberto and J. Tennyson, *Plasma Sources Sci. Technol.* 22, 025001 (2013)



Dissociative-Electron-Attachment

Electron-impact dissociative excitation:

$$e^- + O_2(X^3\Sigma_g^-; v) \rightarrow O_2^-({}^2\Pi_g, {}^2\Pi_u, {}^4\Sigma_u^-, {}^2\Sigma_u^-) \rightarrow e^- + 2 O({}^3P)$$



V. Laporta, R. Celiberto and J. Tennyson, *Phys. Rev.* A91, 012701 (2015). V. Laporta, R. Celiberto and J. Tennyson, *Plasma Sources Sci. Technol.* 22, 025001 (2013)

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CO₂ Plasma based conversion

Recently, worldwide attention has been focused on

increasing emissions of greenhouse gases

decreasing reserves of traditional energy sources

 CO_2

 CH_4

CO₂ Plasma based conversion



CO₂ Plasma based conversion



Highly endothermic carbon dioxide dissociation is one of the key processes affecting the overall reaction efficiency

A promising approach for this conversion involves the non-equilibrium chemistry in the plasma phase by electric discharges

Electron-CO₂ cross sections

Vibrational-excitation

 $e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu) \rightarrow \mathrm{CO}_2^-({}^2\Pi_{\mathrm{u}}) \rightarrow e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu')$

CO₂ SYMMETRIC MODE

MOLPRO (MRCI model, aug-cc-pVQZ basis set, D_{2h}): 3 CORE ORBITALS: $(2A_g + 0B_{3u} + 0B_{2u} + 0B_{1g} + 1B_{1u} + 0B_{2g} + 0B_{3g} + 0A_u)^6$ 12 VALENCE ORBITALS: $(3A_g + 2B_{3u} + 2B_{2u} + 0B_{1g} + 3B_{1u} + 1B_{2g} + 1B_{3g} + 0A_u)^{16}$

CO₂ BENDING/ANTISYMMETRIC MODE

MOLPRO (MRCI model, aug-cc-pVQZ basis set, $C_{2\nu}$): 3 CORE ORBITALS: $(2A_1 + 0B_1 + 1B_2 + 0A_2)^6$ 12 VALENCE ORBITALS: $(5A_1 + 2B_1 + 4B_2 + 1A_2)^{16}$
Letter

Calculated low-energy electron-impact vibrational excitation cross sections for CO₂ molecule

V Laporta^{1,2}, J Tennyson² and R Celiberto^{1,3}

Vibrational-excitation

 $e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu) \rightarrow \mathrm{CO}_2^-({}^2\Pi_{\mathrm{u}}) \rightarrow e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu')$

R-Matrix potential energy curves



Letter

Calculated low-energy electron-impact vibrational excitation cross sections for CO₂ molecule

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Vibrational-excitation

 $e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu) \rightarrow \mathrm{CO}_2^-({}^2\Pi_{\mathrm{u}}) \rightarrow e + \mathrm{CO}_2(\mathrm{X}\,{}^1\Sigma_{\mathrm{g}}^+, \nu')$



Dissociative-attachment for CO₂

Dissociative-excitation for CO₂

 $e + CO_2(X) \rightarrow CO(X) + O^{-}(^2P)$

 $e + CO_2(X) \rightarrow CO(X) + O(^3P) + e$

CO_2/CO_2^-

MOLPRO (MRCI model, aug-cc-pVQZ basis set, $C_{2\nu}$): 5 CORE ORBITALS: (5A₁ + 0B₁ + 0B₂ + 0A₂)¹⁰ 10 VALENCE ORBITALS: (4A₁ + 3B₁ + 3B₂ + 0A₂)¹⁶

In preparation...

Dissociative-attachment for CO₂

Dissociative-excitation for CO₂

 $e + CO_2(X) \rightarrow CO(X) + O^{-}(^2P)$

 $e + CO_2(X) \rightarrow CO(X) + O(^{3}P) + e$



Abdillah Abdoulanziz

In preparation...

Dissociative-attachment for CO₂

Dissociative-excitation for CO₂

 $e + CO_2(X) \rightarrow CO(X) + O^{-}(^2P)$

 $e + CO_2(X) \rightarrow CO(X) + O(^{3}P) + e$



Abdillah Abdoulanziz

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electron-BeH⁺ and -BeD⁺ scattering

$$e + \operatorname{BeH}^{+}(X \ ^{1}\Sigma^{+}; v, j) \to \operatorname{BeH}^{+}(X \ ^{1}\Sigma^{+}; v', j') + e \qquad (VE)$$

$$\to \operatorname{Be}^{*}(n = 1 \dots 12) + \operatorname{H}(^{2}S) \qquad (DR)$$

$$\to \operatorname{Be}^{+}(^{2}S) + \operatorname{H}(^{2}S) + e \qquad (DE)$$

I.T.E.R.

V. Laporta, K. Chakrabarti, et al., Plasma Phys. Control. Fusion **59**, 045008 (2017) S. Niyonzima, V. Laporta, et al., Plasma Sources Sci. Technol. **27**, 025015 (2018)

electron-BeH⁺ and -BeD⁺ scattering



UK-R-Matrix

Potential energy curves and couplings



Figure 1. Potential energy curves (top line) and widths (bottom line) for BeH^{**} resonant states in ${}^{2}\Sigma^{+}$, ${}^{2}\Pi$ and ${}^{2}\Delta$ scattering channels for $j^{+} = 0$. Dashed line refers to potential energy curve for BeH⁺ ground electronic state where the first vibrational levels are marked.

$e + BeH^+$





I.T.E.R.

 $e + BeD^+$





electron-H₂⁺



$$\begin{split} \mathrm{H}_{2}^{+}(N_{i}^{+},v_{i}^{+})+e^{-}(\varepsilon) &\longrightarrow \mathrm{H}+\mathrm{H}\,, & \text{Dissociative Recombination} \\ \mathrm{H}_{2}^{+}(N_{i}^{+},v_{i}^{+})+e^{-}(\varepsilon) &\longrightarrow \mathrm{H}_{2}^{+}(N_{f}^{+},v_{f}^{+})+e^{-}(\varepsilon_{f})\,, & \text{Vibrational Excitation} \\ \mathrm{H}_{2}^{+}(N_{i}^{+},v_{i}^{+})+e^{-}(\varepsilon) &\longrightarrow \mathrm{H}+\mathrm{H}^{+}+e^{-}(\varepsilon_{f})\,. & \text{Dissociative Excitation} \end{split}$$

V. Laporta et al. In preparation...



electron-H₂⁺



electron-H₂⁺





e + ArH⁺ collisions

$e(\varepsilon) + \operatorname{ArH}^+(X^1\Sigma^+, v^+) \to \operatorname{Ar} + \operatorname{H}, (DR)$

A. Abdoulanziz, F. Colboc, D.A. Little, Y. Moulane, J.Zs. Mezei, E. Roueff, J. Tennyson, I.F. Schneider and V. Laporta, MNRAS 479, 2415 (2018)

 $e(\varepsilon) + \operatorname{ArH}^+(X^1\Sigma^+, v^+) \to \operatorname{Ar} + \operatorname{H}, (DR)$



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e + OH collisions

$$e + OH(X^{2}\Pi) \rightarrow e + OH(A^{2}\Pi, a^{4}\Sigma^{-}, 1^{2}\Sigma^{-}), \qquad (EE)$$

$$\rightarrow e + O(^{3}P) + H(^{2}S), \qquad (DE)$$

$$\rightarrow e + O(^{1}D) + H(^{2}S). \qquad (DE)$$

$$\rightarrow O^{-}(^{2}P) + H(^{2}S) \qquad (DA)$$

K. Chakrabarti, V. Laporta, and Jonathan Tennyson, submitted, 2019





Dissociation

Dissociative electron attachment



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Conclusions and perspectives

- Ab-initio molecular Quantum dynamics;
- State specific cross sections and rate coefficients for vibrational excitation, dissociative attachment, dissociative recombination and dissociative excitation processes;
- Non-equilibrium molecular plasmas modelling;
- Extensions to atom-molecule and molecule-molecule scattering in particular for the reaction:

 $Be^+ + H_2 \rightarrow BeH^+ + H$ $Be^+ + H_2 \rightarrow BeH + H^+$

• Data base for $e + H_2$ and $e + H_2^+$ collisions.



Jonathan Tennyson (University College London, UK)



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Arnaud Bultel (CORIA, Rouen, France)



Marco Panesi (University of Illinois, (IL) USA)



Mario Capitelli (CNR, Bari, Italy)



Savino Longo (Università di Bari, Italy)



Carte du monde "électron-molécule"



Carte du monde "électron-molécule"



Carte du monde "électron-molécule"



Non-equilibrium molecular plasmas modelling

- *Non-equilibrium thermodynamics* is an extension of equilibrium thermodynamics it attempts to describe the physical processes in their *time-courses* in continuous detail.
- A central role in non-equilibrium thermodynamics is played by <u>transport equations</u> which are the corresponding to equations of state in equilibrium thermodynamics.
- Transport equations describe the amount of heat, mass, electrical charge... which are transferred per unit time between different systems and different regions of a system as a response to a non-homogeneity in temperature T, molar concentration c, electric potential φ_{e} ...

$$\begin{aligned} q &= -\lambda \nabla T \quad \text{(Fourier's law)}, \\ J &= -D \nabla c \quad \text{(Fick's law)}, \\ I &= \sigma_{\rm e} \nabla \varphi_{\rm e} \quad \text{(Ohm's law)}. \end{aligned}$$



Applications: non-equilibrium dissociation mechanism

- Despite its apparent simplicity, the dissociation rate for a molecule is very difficult to be determined
- Polak shows that experimental dissociation rates cannot be reproduced by direct electron impact dissociation involving vibrational ground state of molecules.
- The "pure-vibrational-mechanism" is the theoretical framework to explain the dissociation rate





M. Capitelli, G. Colonna, V. Laporta et al., Chemical Physics 438, 31 (2014)



M. Capitelli, G. Colonna, V. Laporta et al., Chemical Physics 438, 31 (2014)

Titres universitaires :

2007 – Doctorat en physique, Université de Bari, Italie (séjours à Lyon et au CERN à Genève)

- 2010 Qualification pour maitre de conférences section 29 (France)
- 2017 Qualification pour maitre de conférences section 30 (France)
- 2017 Habilitation scientifique national section 02/B2 et 03/A2 (Italie)
- 2017 Habilitation à Diriger les Recherches, Université du Havre, France
- 2018 Qualification pour professeur section 30 (France)

Parcours professionnel après la thèse :

2006 – 2007 Centre de Physique Théorique, CNRS-École Polytechnique, Palaiseau Cedex, France
2007 – 2009 Dipartimento interateneo di fisica, Università di Bari, Italia
2009 – 2010 Institut de Physique Nucléaire de Lyon, Université Claude Bernard Lyon-1, France (A.T.E.R.)
2010 – 2011 Istituto di Metodologie Inorganiche e dei Plasmi, Consiglio Nazionale delle Ricerche, Bari, Italie
2011 – 2013 Department of Physics and Astronomy, University College London, London (UK); Prof. J. Tennyson
2013 – 2015 Istituto di Nanotecnologia-CNR et DICATECh-Politecnico di Bari (Italie)
Juillet 2013 Ames Research Center, NASA, Moffett Field (CA), USA; Prof. R.L. Jaffe;
2015 – 2015 Ohio Aerospace Institute et Air Force Research Laboratory-WPAFB, Dayton (OH) - USA
2015 – 2016 Istituto di Nanotecnologia-CNR et DICATECh-Politecnico di Bari (Italie)
Juin 2017 Aerospace Engineering Department, University of Illinois at Urbana-Champaign (IL-USA)
De Fevrier 2017 Laboratoire Onde et Milieux Complexes, CNRS-Université du Havre, Le Havre (France)

Application: Electron-vibration relaxation in oxygen plasmas

$$e + O_2(X^3\Sigma_g^-; v, j) \rightleftharpoons e + O_2(X^3\Sigma_g^-; w, j), \quad (VE)$$

- State-to-State vibrational kinetics
- Vibrational relaxation time is comparable to chemical relaxation: vibrational non-equilibrium



V. Laporta, K.L. Heritier and M. Panesi, Chemical Physics 472 (2016) 44–49

Time evolution of **non-equilibrium** vibrational distribution function:



Time evolution of **non-equilibrium** vibrational distribution function:



Time evolution of **non-equilibrium** vibrational distribution function:



Vibrational relaxation time:

