

Cascade Morphologies in BCC Metals

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Introduction

Motivation

- Compare different cascade morphologies in bcc metals.
- Study the effect of morphology on defect production.
- Study the effect on defect clustering/distribution.

Interatomic potentials for MD

- Cr: EAM by Bonny et al [PhilMag 91, 1724 (2011)], ZBL added by N. Juslin for this study.
- Fe: Finnis-Sinclair [PhilMag A 50, 45 (1984)], modified by Calder et al [JNM 207, 25 (1993)].
- Mo: Ackland-Thetford [Philmag A 56, 15 (1987)], modified by Salonen *et al* [JPCM 15, 5845 (2003)].
- W: Ackland-Thetford, modified by Juslin *et al* [JNM **432**, 61 (2013)], denoted as ATJ potential.
- W-björkas: Derlet-Dudarev [PRB **76**, 054107 (2007)], modified by Björkas *et al* [NIMB **267**, 3204 (2009], denoted as DDB potential.

Simulations were performed with LAMMPS code at 100 K (Fe) and 300 K (Cr, Mo, W). Electronic loss was excluded.

Defect was determined based on occupancy of Wigner-Seitz cell.

More information of the setup: JPCM 27, 225402 (2015).



Observation of Different Cascade Regimes



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- A kink in the N_F curve is evident in W.
- We propose that the kink also exists in Cr, Fe, and Mo.
- Hence, we perform a bilinear fit and determine the inflection point (transition energy).
- The transition energy, μ , marks the onset of a different power-law regime.
- Depending on one's view, there could be a transition at the low-energy end: before and after "cascade-like" behavior is fully developed, as suggested in Fe [JNM 251, 49 (1997)] and W [JNM 462, 329 (2015)].

Displacement Threshold Energy, E_d



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$$E_{\rm d} = \sum_i E_{d,i} P_i \Delta \Omega_i / \sum_i P_i \Delta \Omega_i,$$

$$\Delta\Omega_i(h,k,l) = 2\pi \left(1 - \frac{\sqrt{h^2 + k^2 + l^2}}{\sqrt{h^2 + k^2 + l^2 + D^2/4}} \right),$$

Interatomic potential	$E_{\rm d}~({\rm eV})$
Cr	34.4 ± 1.5
Fe	41.8 ± 1.6
Mo	78.7 ± 2.9
W	122.6 ± 4.4





Normalization of Cascade Energy Scale



- To find a possible scaling of μ , we normalize the cascade energy with the displacement threshold energy, E_d .
- Indeed, the μ of Cr, Fe, Mo, and W fall in a narrow range of the normalized energy scale, around $\mu/E_d\approx 300.$
- μ also marks the onset of the formation of large SIA clusters in W and to a lesser extent in Mo, Fe, and Cr.

Cascade Morphology Comparison



Interacting multiple supersonic shocks facilitate the formation of the large SIA clusters and vac clusters, was first suggested by Calder *et al* in Fe with 2 Bi PKAs [PhilMag **90**, 863 (2010)].

Effect of Cascade Morphology on Size Distribution of Defect Clusters in W



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Slope change \rightarrow associated with different cluster formation mechanism of the underlying cascade morphology.

Effect of Cascade Morphology on Size Distribution of Defect Clusters in W



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Large clusters form more frequently in an interacting multiple-shock cascade than expected from a single-shock cascade.

Figure 2. Size–frequency distributions of defects in W derived from MD simulations of cascades initiated by PKAs with various energies. Results are averaged over 20, 36 and 10 cascades for 100, 150 and 200 keV PKAs, respectively. Note the different slopes of scaling laws for SIAs and vacancies. For clarity the vacancy plot is shifted upwards with respect to the SIA plot. The scale of the axes is the same for both cases.

A. E. Sand et al, MatResLett 5, 357 (2017).



Additional BCC Test Systems

Interatomic potential	$E_{\rm d}~({\rm eV})$	Note	
Cr	34.4 ± 1.5	As is.	
Fe	41.8 ± 1.6	As is.	
Mo	78.7 ± 2.9	As is.	
W	122.6 ± 4.4	As is.	
Cr	33.8 ± 1.4	With W mass.	(Cr/W-mass)
Cr	18.9 ± 1.1	With W PKAmass.	(Cr/W-PKAmass)
W	123.6 ± 4.4	With Cr mass.	(W/Cr-mass)
W	170.9 ± 6.7	With Cr PKAmass.	(W/Cr-PKAmass)
W-björkas	98.0 ± 3.7	As is.	

- PKAmass \clubsuit E_d \clubsuit , and vice versa, related to momentum transfer.
- All-atom mass has negligible effect on E_d, due,
 to similar collision kinematics.



Transition Energy Model

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- μ is found to be reasonably proportional to E_d:
 - \checkmark The transition energy can be estimated based on only one parameter, i.e. E_d .
 - ✓ The effect of different mass, PKA mass, and the force field on μ is embedded in their effect on E_d.



Linear vs Bilinear Fit of N_F Data

System	Two-line fit $\delta_{\rm rms}(\%)$	One-line fit $\delta_{\rm rms}(\%)$	
Cr	5.1	10.2	-
Fe	5.9	18.8	Γ
Мо	5.0	9.3	$1 \frac{N}{N} \left(y_{1} - y_{2} \right)^{2}$
W	3.3	19.9	$\delta_{\rm rms} = 100 \left\lfloor \frac{1}{2} \sum \left(\frac{y_i - y_{i,\rm fit}}{y_i} \right) \right\rfloor$
Cr/W-mass	5.8	12.4	$N \stackrel{\text{orms}}{\longrightarrow} N \stackrel{\text{res}}{\longrightarrow} V_{i,\text{fit}}$
Cr/W-PKAmass	5.3	18.4	\mathbf{N} $i=1$ \mathbf{C} \mathbf{C} \mathbf{C}
W/Cr-mass	5.6	24.2	
W/Cr-PKAmass	5.5	11.8	
W-björkas	6.3	34.0	

• The two-line (bilinear) fit performs significantly better than the one-line (linear) fit.

Early Studies of Sub-Cascade with Binary Collision Approximation (BCA)



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Heinisch et al, JNM **179**, 893 (1991)

10-1 Element Reduced Cascade Reduced break-up break-up energy factor^{a)} recoil damage (eV^{-1}) energy energy Vacancy Density (a²³ (keV) 2.34×10^{-5} **A**1 2.5 0.052 Cu 3.05×10^{-6} 20 0.049 1.28×10^{-6} 50 0.055 10⁻³, Ag FCC Cascade Break-up Energy 4.18×10^{-7} Au 200 0.071 Fe 4.00×10^{-6} 50 0.15 Mo 1.80×10^{-6} 130 0.19 5.10×10^{-7} W 300 0.13 10-4 -3 10⁻² 10⁻¹ Reduced Energy = Recoil E / V(r = screening radius) -3 10

- BCA shows that cascade transition occurs both in bcc and fcc metals and over a narrow range of reduced cascade energy.
- Only non-interacting cascades (sub-cascades) were predicted by using BCA (many body collisions are ignored). MD shows that W exhibits interacting cascade (no sub-cascades).
- In W, will the next transition occur? i.e. interacting cascade \rightarrow non-interacting cascade.

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300-keV Cascade in Tungsten

E. Zarkadoula *et al*, JPCM **27**, 135401 (2015). 2T-model = electronic stopping (friction model) + electron-phonon coupling E_{PKA} = 300 keV with DDB potential.



Non-interacting cascades are still not observed at 300 keV.



Electronic Stopping Effect in Tungsten

A. Sand et al, EPL 103, 46003 (2013).

A. Sand et al, JNM 455, 207 (2014).

Electronic stopping (ES) was modeled as a friction term when atom's kinetic energy $> T_c$.

 E_{PKA} = 150 keV, simulation at 0 K.



 Compared to the friction model, ignoring ES overestimates N_F even though still reasonably within the standard errors.

Electronic Loss Effect on Size Distribution of Defect Clusters in W



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Large cluster probability is smaller with electronic loss.

A. E. Sand *et al*, EPL **115**, 36001 (2016).

A. Sand *et al*, EPL **103**, 46003 (2013). A. Sand *et al*, JNM **455**, 207 (2014).



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Path of energetic atoms from 0 fs (red) to 200 fs (blue) showing interconnected regions, consistent with our results.

Ignoring ES increases potential energy \rightarrow energy density \rightarrow larger defect clusters. Nevertheless, friction model also predicts interacting multiple-shock morphology.

Electron-Phonon Coupling Effect Tungsten

E. Zarkadoula *et al*, JPCM **27**, 135401 (2015). 2T-model = electronic stopping (friction model) + electron-phonon coupling E_{PKA} = 300 keV with DDB potential.



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Through out simulation, $T_{atom} > T_{electron} \rightarrow electrons$ act as additional heat sinks. N_{def} with 2T model < with friction model (note that the amount of net damage energy is not the same, hence the net effect of the 2T model is difficult to infer).



Summary

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- Three cascade morphologies are observed:
 - Single-shock cascade (before sub-cascade).
 - Separate (non-interacting) multiple-shock (sub-cascade in Cr, Fe, Mo).
 - Interacting multiple-shock (connected sub-cascade in W) → formation of large clusters.

A transition energy separates different energy regimes of cascade morphology:

- different power-law in defect production and cluster size distributions.
- Hence data should be fit with a bilinear curve.

In simple bcc metals, the transition energy occurs at damage energy approximately 300E_d.